



Recent advances on Dynamics at Météo-France.

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Context

Main features of the current dynamical kernel of IFS/ARPEGE Global Models and ALADIN/HIRLAM/ALARO/HARMONIE/AROME (LAMs)

- Access by a switch to Fully-compressible (EE) or Hydrostatic (HPE) governing equations cast in time-dependent hybrid mass-based terrain-following η -coordinate.
- Horizontal Spectral (SP) transform discretisation.
- Semi-implicit (SI) time-stepping techniques.
- Semi-Lagrangian (SL) transport scheme.

Need for changes

• In the prospect of future HPC architectures, current status of the dynamical core need to be revisited for more flexible, efficient and scalable numerical algorithms.

Overview

- SWITCH TO (MODERN) SOUND-PROOF NH EQUATIONS
 - A Quasi-elastic system cast in mass-based coordinate
- 2 ABANDONING SPECTRAL TECHNIQUE ?
 - Accurate grid-point derivatives and diffusion
 - Grid-point SI mass-based Solvers
- AN ALTERNATIVE TO SEMI-IMPLICIT TIME-SCHEMES ?
 - RK-IMEX HEVI schemes

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Motivations

- Suppress high-frequency vertically-propagating acoustic wave at their source
 Potential benefit in term of stability, and for initialization procedure.
- Design an approximate set of sound-proof NH equations in mass-based coordinate, accurate at both small- and large- scales, by exploiting Arakawa and Konor (2009) idea together with Laprise (1992) formalism.
- Formulate "blended" numerical model, where different model configurations (soundproof-to-fully compressible) are seamlessly accessed by straightforward switching within a single numerical framework.

Definition of pseudo-hydrostatic reference-state

Defining $(\pi, \widetilde{\rho}, \widetilde{T})$, the pseudo-hydrostatic thermodynamic variables by

$$\frac{\partial \pi}{\partial z} = -\widetilde{\rho} g$$

$$\widetilde{\rho} = p_{00}^{\kappa} \frac{\pi^{1-\kappa}}{R\theta} \exp[\xi(1-\kappa)q]$$

$$\widetilde{T} = \frac{\pi}{R\widetilde{\rho}} \exp[\xi q]$$

where $\kappa = R/C_p$, and q denotes the pressure departure from the pseudo-hydrostatic pressure π , so that the true thermodynamic variables (p, ρ, T) are entirely determined by

$$egin{array}{lcl}
ho &=& \pi \exp[q] \
ho &=& \widetilde{
ho} \exp[(1-\xi)(1-\kappa)q] \ T &=& \widetilde{T} \exp[(1-\xi)\kappa q] \end{array}$$

 ξ : binary switch for Quasi-elastic soundproof approximation.

Basic underlying idea

True density ρ is replaced by the pseudo-density $\widetilde{\rho}(\theta, \pi, q)$ is the mass-continuity equation:

$$\frac{D\widetilde{\rho}}{Dt} = \left(\frac{\partial\widetilde{\rho}}{\partial\theta}\right)_{\pi,q} \frac{D\theta}{Dt} + \left(\frac{\partial\widetilde{\rho}}{\partial\pi}\right)_{\theta,q} \frac{D\pi}{Dt} + \left(\frac{\partial\widetilde{\rho}}{\partial q}\right)_{\theta,\pi} \frac{Dq}{Dt} = -\widetilde{\rho} \, \mathbb{D}_{3}$$

$$+ \frac{1}{2} \left(\frac{\partial\widetilde{\rho}}{\partial\theta}\right)_{\theta,\pi} \frac{Dq}{Dt} + \frac{1}{2} \left(\frac{\partial\widetilde{\rho}}{\partial\theta}\right)_{\theta,\pi} \frac{Dq}{Dt} = -\widetilde{\rho} \, \mathbb{D}_{3}$$

where

$$\left(rac{\partial \widetilde{
ho}}{\partial oldsymbol{q}}
ight)_{\! heta,\pi} = oldsymbol{\xi} (1-\kappa)\widetilde{
ho}$$

 \mathbb{D}_3 : three-dimensional divergence of wind.

QE approximation $(\xi = 0)$

- Minimal condition for filtering vertically propagating acoustic waves,
- NH compressibility of the fluid is suppressed,
- Hydrostatic compressibility is maintained for good accuracy at large-scales.

Prognostic equations cast in mass-based η -coordinate with $\overline{\pi = A(\eta) + B(\eta)\pi_s}$

$$\frac{DV}{Dt} = -\frac{\widetilde{\rho}}{\rho} \left[\frac{1}{m} \frac{\partial p}{\partial \eta} \nabla \phi + \frac{1}{\widetilde{\rho}} \nabla p \right] + \mathcal{V}$$

$$\gamma \frac{Dw}{Dt} = g \frac{\widetilde{\rho}}{\rho} \left[\frac{1}{m} \frac{\partial (p - \pi)}{\partial \eta} - \frac{\rho - \widetilde{\rho}}{\widetilde{\rho}} \right] + \gamma \mathcal{W}$$

$$\frac{D\widetilde{T}}{Dt} = \frac{R\widetilde{T}}{C_p} \left(\frac{\dot{\pi}}{\pi} + \xi \frac{Dq}{Dt} \right) + \frac{Q\widetilde{T}}{C_p T}$$

$$\xi \frac{Dq}{Dt} = -\frac{\dot{\pi}}{\pi} - \frac{C_p}{C_v} \mathbb{D}_3 + \frac{Q}{C_v T}$$

$$\frac{\partial \pi_s}{\partial t} = -\int_0^1 \nabla \cdot (mV) \, d\eta'$$

 γ : binary switch for Hydrostatic approximation.

Diagnostic equations :

$$m = \frac{\partial \pi}{\partial \eta}$$

$$\nabla \phi = \nabla \phi_{s} + \int_{\eta}^{1} \nabla \left(m \frac{R \widetilde{T}}{\pi} e^{-\xi q} \right) d\eta'$$

$$\frac{\dot{\pi}}{\pi} = V \cdot \frac{\nabla \pi}{\pi} - \frac{1}{\pi} \int_{0}^{\eta} \nabla \cdot (mV) d\eta'$$

$$m \dot{\eta} = B(\eta) \int_{0}^{1} \nabla \cdot (mV) d\eta' - \int_{0}^{\eta} \nabla \cdot (mV) d\eta'$$

$$\mathbb{D}_{3} = \nabla \cdot V + \frac{\widetilde{\rho}}{m} \nabla \phi \cdot \frac{\partial V}{\partial \eta} - g \frac{\widetilde{\rho}}{m} \frac{\partial w}{\partial \eta}$$

Different model configurations: "Switchable" equation set

- **① EE** case $(\xi, \gamma) = (1, 1) : [V, w, \widetilde{T}, q, \pi_s]$ are prognostic variables.
- **4 PPE case** $(\xi, \gamma) = (1, 0) : [V, \widetilde{T}, \pi_s]$ are prognostic variables, $p = \pi \Rightarrow q = 0$ and pressure Eq. is used to diagnose w.
- **QE** case $(\xi, \gamma) = (0, 1) : [V, w, \widetilde{T}, \pi_s]$ are prognostic variables, and q is diagnosed from the pressure Eq., which becomes a divergence constraint relationship.

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Two parallel approaches explored at ECMWF + M-F

- FVM (Smolarkiewicz P., Kuehnlein C.,...) effort focused on: local coordinate, unstructured grid, conservation.
- GlobSWAR (Bénard, Caluwaerts,...) effort focused on: global spherical coord., structured grid, high-order in space also trying to keep most of the existing architecture of IFS/ARPEGE models

"Keep most existing architecture" means:

- Global Spherical coordinate
- Reduced Lat-Lon grid (semi-structured)
- Unstaggered grid (A-grid)

Spherical coord. + Reduced A-grid: not the "easiest" choice!

Two psychological obstacles must be overcome:

- Well-known "pole problem(s)" of spherical coordinates, especially with reduced lat-lon grids
- Well-known gravity wave distortion problem of A-grid

But some advantages also:

- semi-structured grid is a potential advantage (efficiency, easier access to high-order)
- A-grid has the best balance of DoFs ever.
- A-grid is much better for transport/advection than C-grid

⇒ We build a Proof-of-Concept SW model: GlobSWAR

Pole problem – special care required :

- very accurate zonal derivatives near poles
 (e.g. Fourier representation in a small cap)
- very careful definition of BCs at poles for meridional derivatives
- some continuous formulations more convenient than others (e.g. curl form mom. eqs, flux form cont. eq, term grouping...)
- Zonal averages $\langle u \rangle, \langle v \rangle$ must have specific BCs for meridional derivatives (for stability reasons)

A-Grid

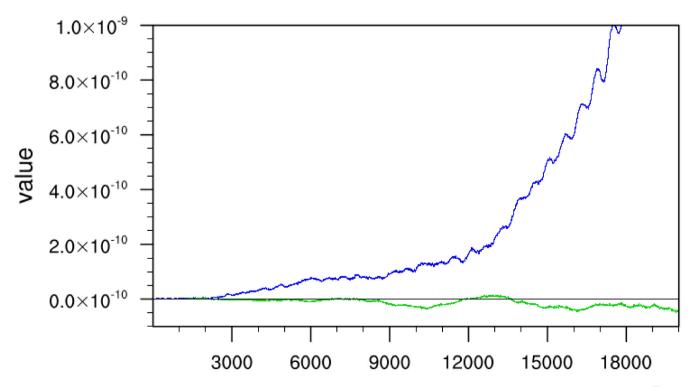
- needs a good filter beyond $3\Delta x$ wave (required for anti-alias of ψ^2 in any case)
- no special problem encountered with A-grid

GlobSWAR: PSBR test-case

• Demonstrate ability to keep simple stationary parallel flows, actually stationary without any damping/diffusion \rightarrow tilted Solid-Body rotation test-case for ($\alpha = 0^{\circ}, 45^{\circ}, 90^{\circ}$).

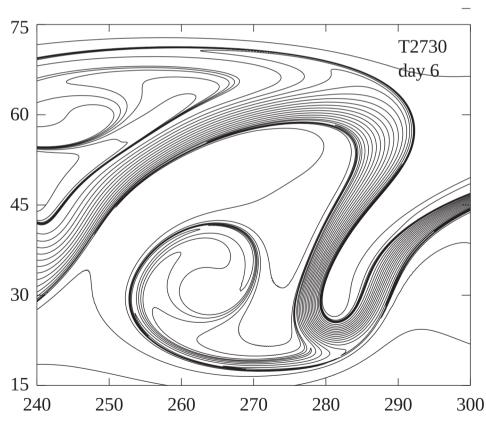
Blue : BC problem for $\langle u \rangle, \langle v \rangle$ not solved.

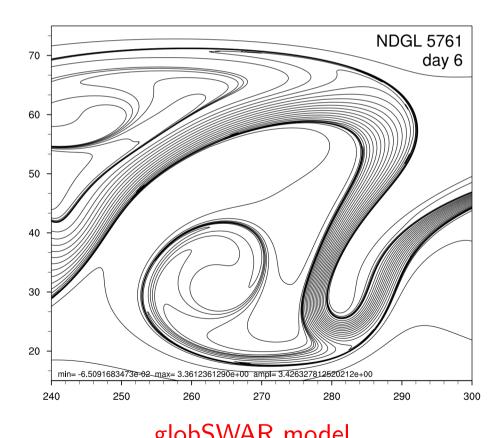
Green: BC problem solved



GlobSWAR: Scott et al. (2016) baroclinic test-case

 Demonstrate ability to perform demanding cases at high resolution (at 3.5 km grid mesh). Barotropric instability pushed near end of determinism Reference solutions = models BOB (Columbia Univ.), and FV3 (GFDL)





BOB model

Grid-point SI mass-based solvers

State of art of current constant-coefficient Spectral SI mass-based NH EE solver

- Horizontal and vertical divergence resp. (D, d) are used as auxiliary prognostic variables in implicit part of the SI scheme.
- At the price of the fulfilment of two vertical constraints, the full linear implicit system is eventually reduced to

$$\begin{bmatrix} I - \delta t^2 \mathbf{B}_{\mathsf{d}}^* \Delta \end{bmatrix} d^+ = d^{\bullet}$$
$$\begin{bmatrix} I - \delta t^2 c_*^2 \Delta \end{bmatrix} D^+ = D^{\bullet} + \delta t^2 \mathbf{G}_{\mathsf{d}}^* d^+$$

- \mathbf{B}_{d}^{*} is a definite positive constant-coefficient vertical matrix, c_{*} is the constant speed of sound.
- d^+ and D^+ are obtained by solving successively **two** 2D horizontal Helmholtz-like equations for each vertical eigenmodes of \mathbf{B}_d^* , enjoying the fact that the horizontal laplacian operator Δ is **diagonal** in spectral space.

Grid-point SI mass-based solvers

A constant-coefficient Grid-point SI mass-based NH solver

• Elimination must be performed in a different manner, relaxing the less restrictive of the two vertical constraints:

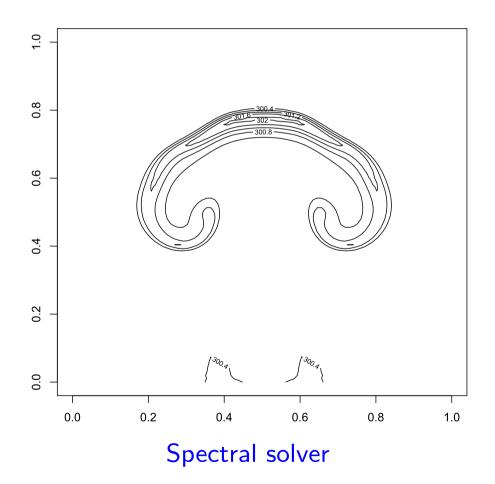
$$egin{aligned} \left[\mathbf{I} - \delta t^2 \mathbf{B}_{\mathsf{D}}^*
abla^2
ight] D^+ &= D^{ullet ullet} \ \mathbf{H}_{\mathsf{v}}^* q^+ &= Q^{ullet} + \delta t^2 \mathbf{S}_{\mathsf{D}}^* D^+ \end{aligned}$$

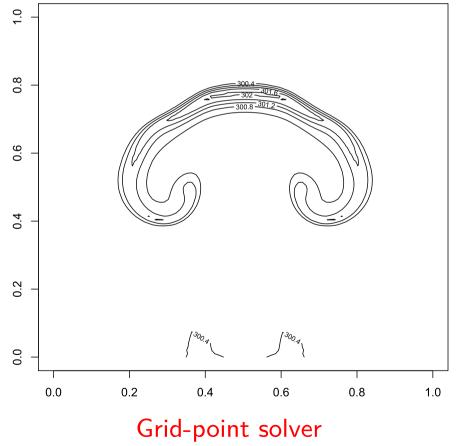
- $\mathbf{B}_{\mathrm{D}}^{*}$ is a newly defined constant definite positive vertical matrix ensuring that $\lambda(\mathbf{B}_{\mathrm{D}}^{*}) = \lambda(\mathbf{B}_{\mathrm{d}}^{*})$, and $\mathbf{H}_{\mathrm{v}}^{*}$ is invertible tridiagonal constant vertical matrix.
- A single 2D horizontal Helmholtz-like equation for each vertical eigenmode of \mathbf{B}_{D}^{*} is solved for projection of D^{+} using iterative of Krylov solvers : GMRES or GCR(k).
- A simple tridiagonal **vertical inversion** to get q^+ .

Grid-point SI mass-based solvers

Test in a vertical Slice $(x-\eta)$ version of AROME model: Rising bubble test-case

 GMRES grid-point Krylov solver using 6th-order finite differences derivatives is compared to current Spectral solver.





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Basic principle of HEVI time-stepping approaches for solving EE equations

- ullet Implicit time treatment of the vertical adjustment terms responsible for vertically propagating high-frequency acoustic modes \Rightarrow 1D vertical Helmholtz-like problems to solve.
- Explicit time treatment of the remaining processes \Rightarrow stability is mostly control by the horizontal wave Courant (CFL) number.

"Our" Desirable properties for HEVI schemes

- No time filtering applied \Rightarrow Multi-steps time-schemes are discarded \Rightarrow Multi-stages time-stepping is favoured.
- No artificial divergence Damping either \Rightarrow Classical time-splitting approach is not retained \Rightarrow "monolithic" temporal treatment of all processes, i.e using a single time-step.
- ⇒ Recent RK-IMEX HEVI schemes appear to be good candidates

Classical RK-IMEX HEVI schemes with double Butcher's tableaux

$$\partial_t X = \mathscr{M}(X) = \mathscr{E}(X) + \mathscr{I}(X).$$
"explicit"

Runge-Kutta scheme for each part:

$$X^{(j)} - \underset{j=1}{\textit{ajj}} \delta t \mathscr{I}(X^{(j)}) = X^t + \delta t \sum_{k=1}^{j-1} \tilde{a}_{jk} \mathscr{E}(X^{(k)}) + \delta t \sum_{k=1}^{j-1} a_{jk} \mathscr{I}(X^{(k)})$$

$$X^{t+\delta t} = X^t + \delta t \sum_{j=1}^{\nu} \tilde{b}_{j} \mathscr{E}(X^{(j)}) + \delta t \sum_{j=1}^{\nu} b_{j} \mathscr{I}(X^{(j)})$$

Coefficients are resumed in a double Butcher's tableau : $\{\tilde{c}, \tilde{\mathcal{A}}, \tilde{b}\}$ and $\{c, \mathcal{A}, b\}$ (where $\tilde{\mathcal{A}} = (\tilde{a}_{ji})_{1 < j, i < \nu}$ et $\mathcal{A} = (a_{ji})_{1 < j, i < \nu}$).

HEVI schemes constraints

- Vertical adjustment terms must to be deat with implicitly and all advection terms have to be treated explicitly.
- The horizontal momentum equations and pressure (mass continuity) equation may be subjected to Foreward/backward time treament in the same spirit as Mesinger [1977]: Wind-Forward and Pressure-Backward (UFPreB) (Lock *et al.* [2014]).
- Pushing this idea further, we exploit the fact that Pressure-Forward and Wind-Backward (UBPreF) is also a possible option to elaborate more stable RK-IMEX HEVI scheme in presence of advection, Colavolpe et al. [2017].

A proposed "Mixed" four Butcher's tableaux RK-IMEX HEVI scheme

$$\partial_t X = \underbrace{\mathcal{A}(X)}_{\text{hor.grad}(p)} + \underbrace{\mathcal{J}(X)}_{\text{hor.grad}(p)} + \underbrace{\mathcal{U}(X)}_{\text{hort.div}(\mathbf{v})} + \underbrace{\mathcal{J}(X)}_{\text{hort.div}(\mathbf{v})} + \underbrace{\mathcal{J}(X)}_{\text{hort.div}(\mathbf{v})$$

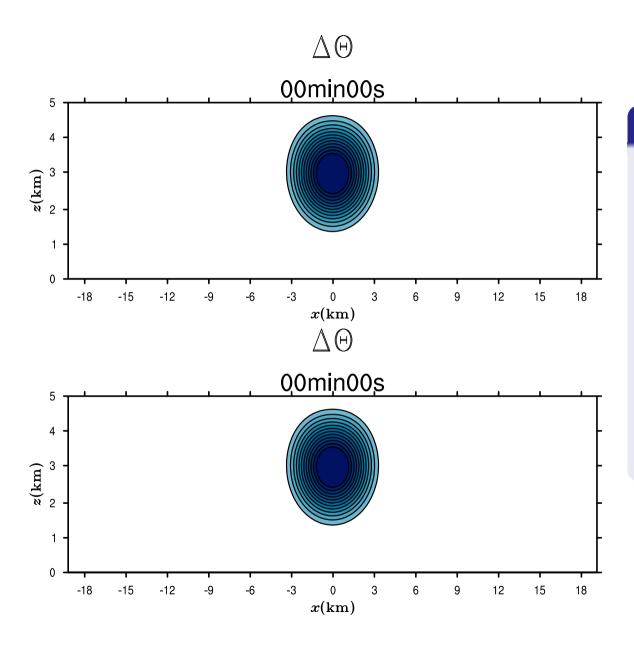
Hence

$$\frac{X^{(j)} - X^{t}}{\delta t} = \sum_{k=1}^{j-1} \tilde{a}_{jk} \mathscr{A}(X^{(k)}) + \sum_{k=1}^{j} a_{jk} \mathscr{I}(X^{(k)}) + \sum_{k=1}^{j} a_{jk}^{u} \mathscr{U}(X^{(k)}) + \sum_{k=1}^{j} a_{jk}^{p} \mathscr{P}(X^{(k)})$$

$$\frac{X^{t+\delta t} - X^{t}}{\delta t} = \sum_{i=1}^{\nu} \tilde{b}_{j} \mathscr{A}(X^{(j)}) + \sum_{i=1}^{\nu} b_{j} \mathscr{I}(X^{(j)}) + \sum_{i=1}^{\nu} b_{j}^{u} \mathscr{U}(X^{(j)}) + \sum_{i=1}^{\nu} b_{j}^{p} \mathscr{P}(X^{(j)})$$

- Coefficients of the four Butcher's tableau $\{\tilde{c}, \tilde{\mathcal{A}}, \tilde{b}\}$, $\{c, \mathcal{A}, b\}$, $\{c^u, \mathcal{A}^u, b^u\}$ and $\{c^p, \mathcal{A}^p, b^p\}$ must satisfied Colavolpe *et al.* [2017] conditions.
- Alternative "Foreward/Backward" treatment of "horizontal adjustment terms" : pressure gradient/wind divergence.
- No computational overcost, second-order accuracy, storage and number of vertical inversions.

Abandoning SI time-schemes



Straka et al. (1993) test-case

$$\Delta x == 75 \,\mathrm{m}, \ \Delta z \approx 75 \,\mathrm{m},$$

 $\Delta t = 0.125 \,\mathrm{s} \Rightarrow C^* \approx 1.8$

8-th order horizontal FD, 2-nd order vertical FD.

- **Top** : Classical UJ13(1,3,2) double butcher's tableaux.
- **Bottom**: New UJ13(1,3,2) mixed four butcher's tableaux.

On-going Works

- Implementation of the QE dynamical system as a switch option in the full 3D common code is in progress (almost done). Testing and validations will be the next moves.
- A semi-implicit version of GlobSWAR model is first envisaged, before extending to multi-level GlobHPAR model.
- We are now ready to introduce a more complex linear SI operator including the variable-coefficient in order to deal with steep orography. This study will be pursued in the Vertical slice version of AROME using Krylov Grid-point Solver approach.
- 3D prototype of mass-based EE model integrating over a Cartesian domain through our best RK-IMEX HEVI time-scheme should be available soon.