

# Recent numerics developments in the COSMO model and an outlook for the ICON model

41<sup>st</sup> EWGLAM / 26<sup>th</sup> SRNWP-meeting, Sofia, Bulgaria  
30 Sept. – 03 Oct. 2019

**Michael Baldauf** (DWD), **Damian Wojcik** (IMGW)

## Consortium in transition

The current **COSMO model** (with the ‚Runge-Kutta‘ dyn.core) is slowly phased out during the next years in the COSMO consortium.

- DWD plans to replace COSMO-D2 by **ICON-D2** in Q4/2020
- ...
- MeteoCH plans the replacement ~2023 (← adaptations to GPU computers!)

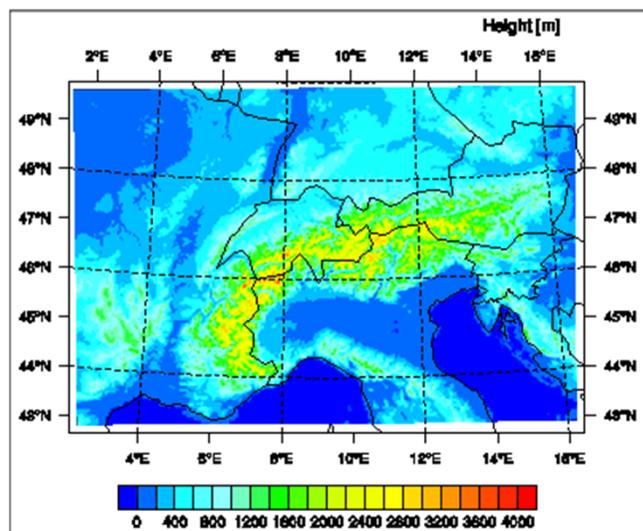
This migration is prepared by all COSMO partners in the priority project ‚*Transition of COSMO to ICON-LAM (C2I)*‘; project leader: Daniel Rieger (DWD)

Therefore, no further development work at the dyn. core will be done from now on (exception: higher-order scheme by Univ. Cottbus, A. Will)

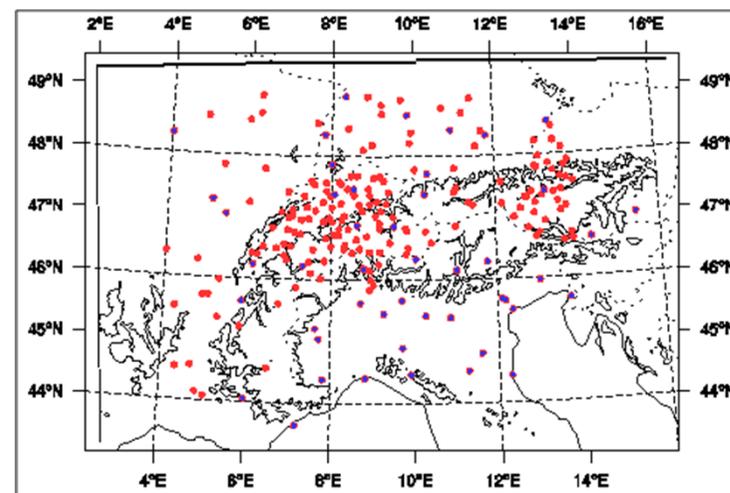
However, **COSMO-EULAG** will be further developed and probably will go into operations at IMGW (Poland) (currently pre-operational)

Investigations:

1. **TKE advection:** replacement of the Bott scheme by MPDATA-A
2. Selection of an optimal MPDATA advection version



Topographical map of the domain



Station network for surface verification

Setup of experiments:

- Operational COSMO-2 domain used by Meteo-Swiss, 60 vertical levels
- Entire June 2013, 48-hour forecasts, verification by VERSUS software
- Numerical and Smagorinsky diffusion are *turned off* for COSMO-EULAG

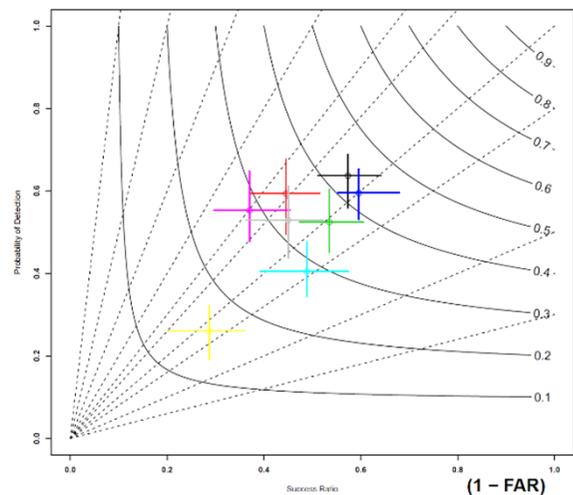
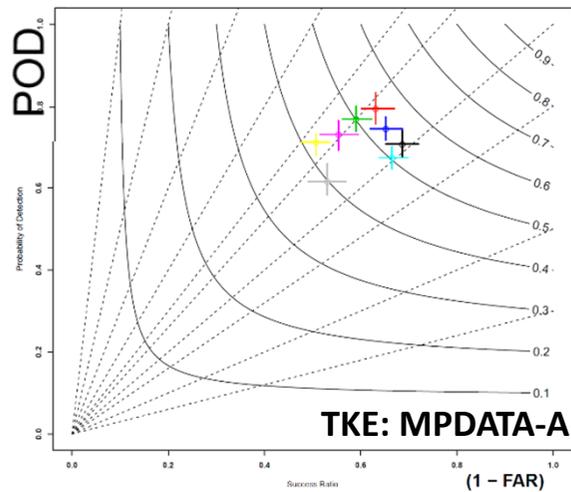
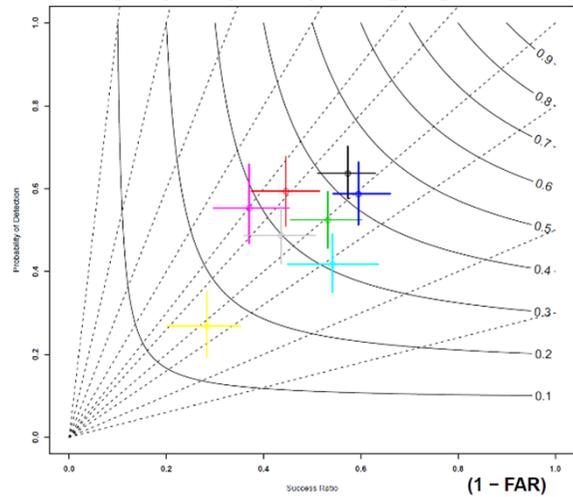
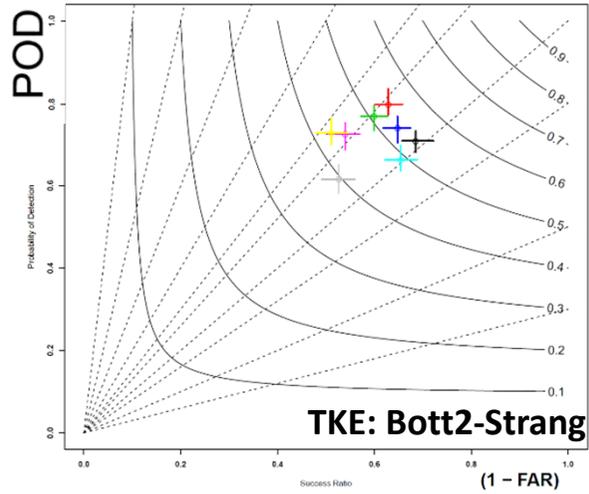
# Replacement of the Bott TKE adv. scheme by the MPDATA-A scheme

The verification scores of COSMO-EULAG do not alter significantly with that change.

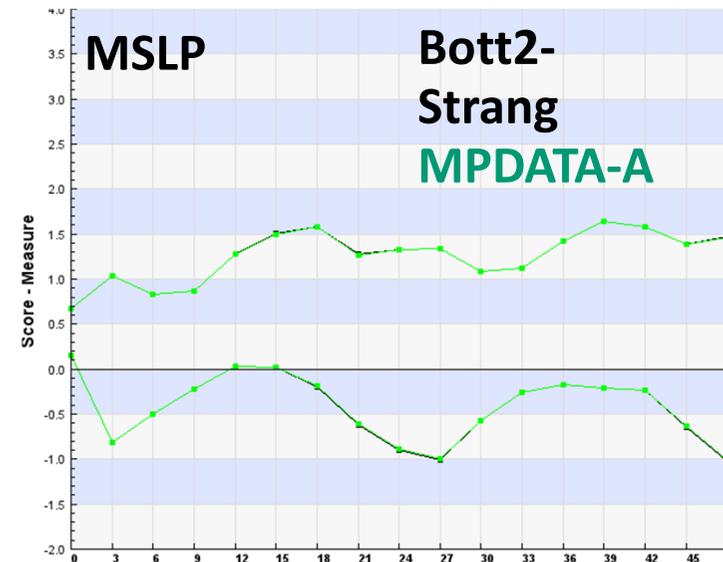
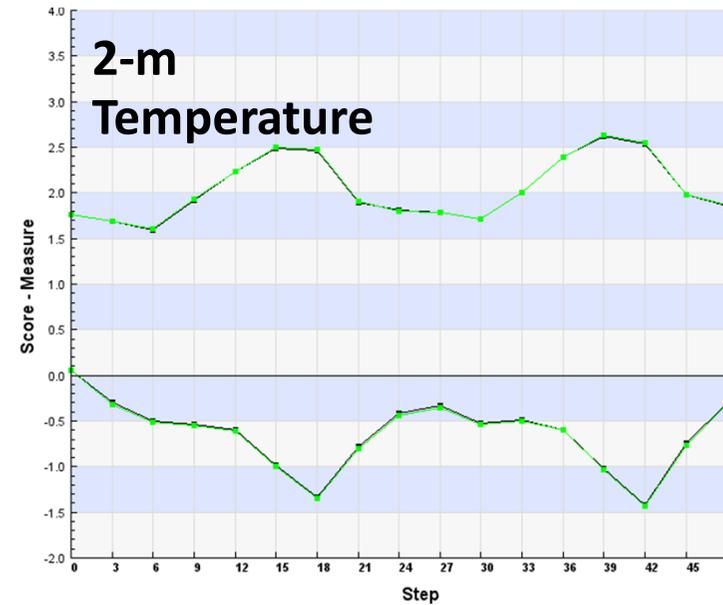
1 mm and more

8mm and more

○ step 6 ○ step 12 ○ step 18 ○ step 24 ○ step 30 ○ step 36 ○ step 42 ○ step 48

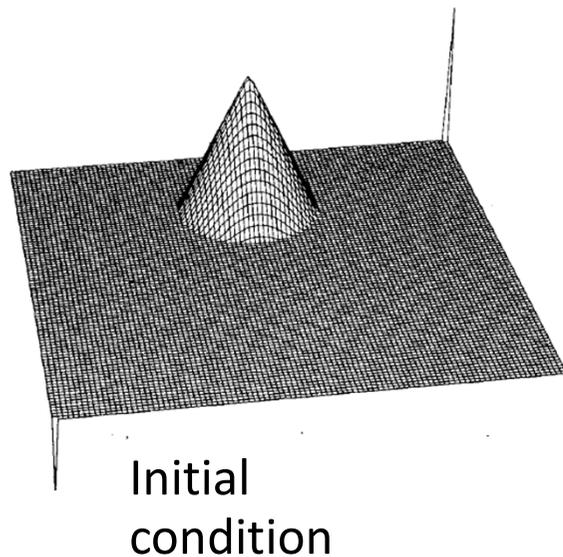
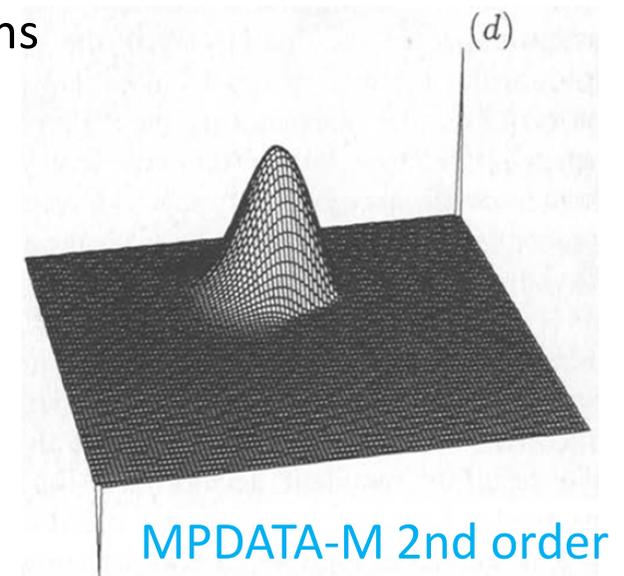
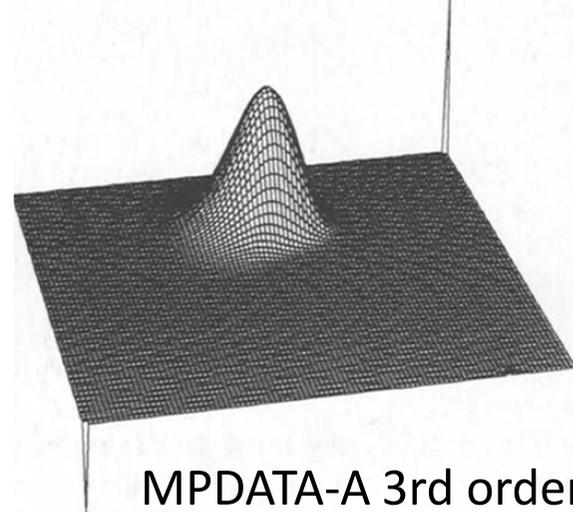
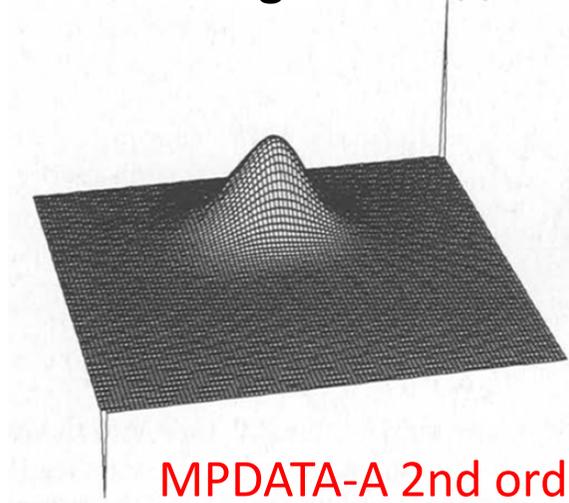


Reliability diagrams for precipitation



# Selection of an optimal adv. sch. (MPDATA-A vs. MPDATA-M)

The Rotating Cone (a) Test Case – results after 6 (b) revolutions



	MPDATA-A 2 <sup>nd</sup> order	MPDATA-M 2 <sup>nd</sup> order
Advected field	$\psi$	$\Psi+c, c \rightarrow \infty$
Accuracy	Lower	Higher
Diffusion	More	Less

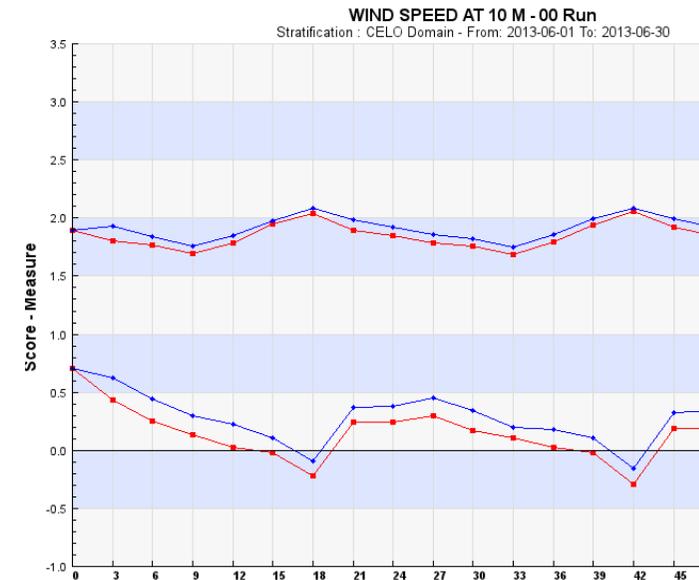
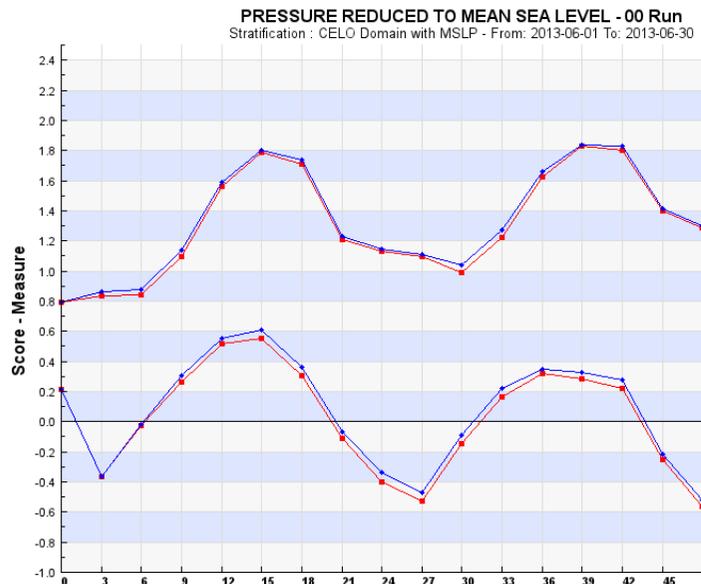
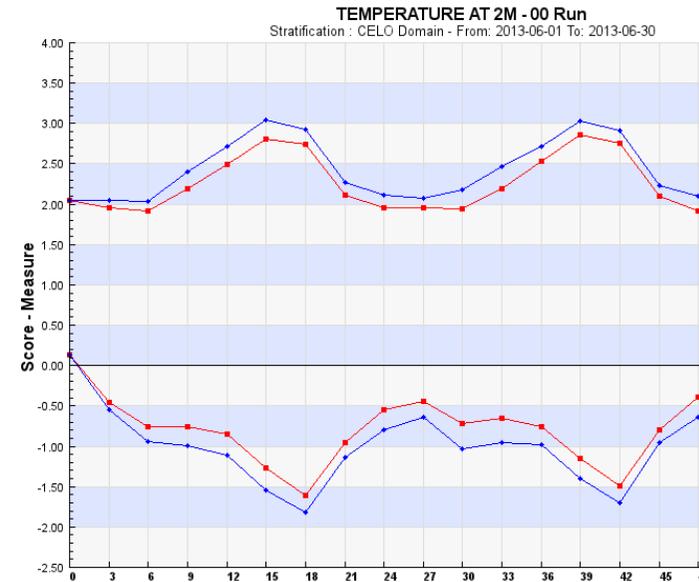
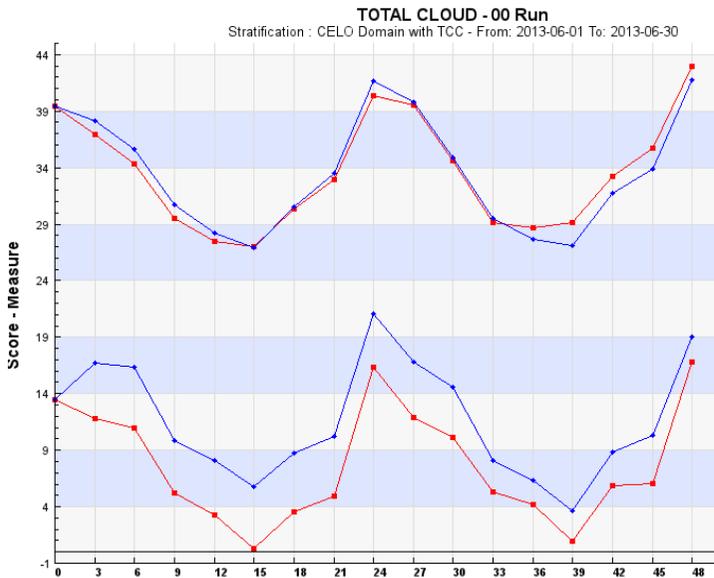
MPDATA = upwind advection with iterative improvement by 'anti-diffusive' fluxes (nonlinear)

Smolarkiewicz and Clark (JCP, 1986),  
Smolarkiewicz and Grabowski (JCP, 1990).

Damian Wojcik (IMGW)



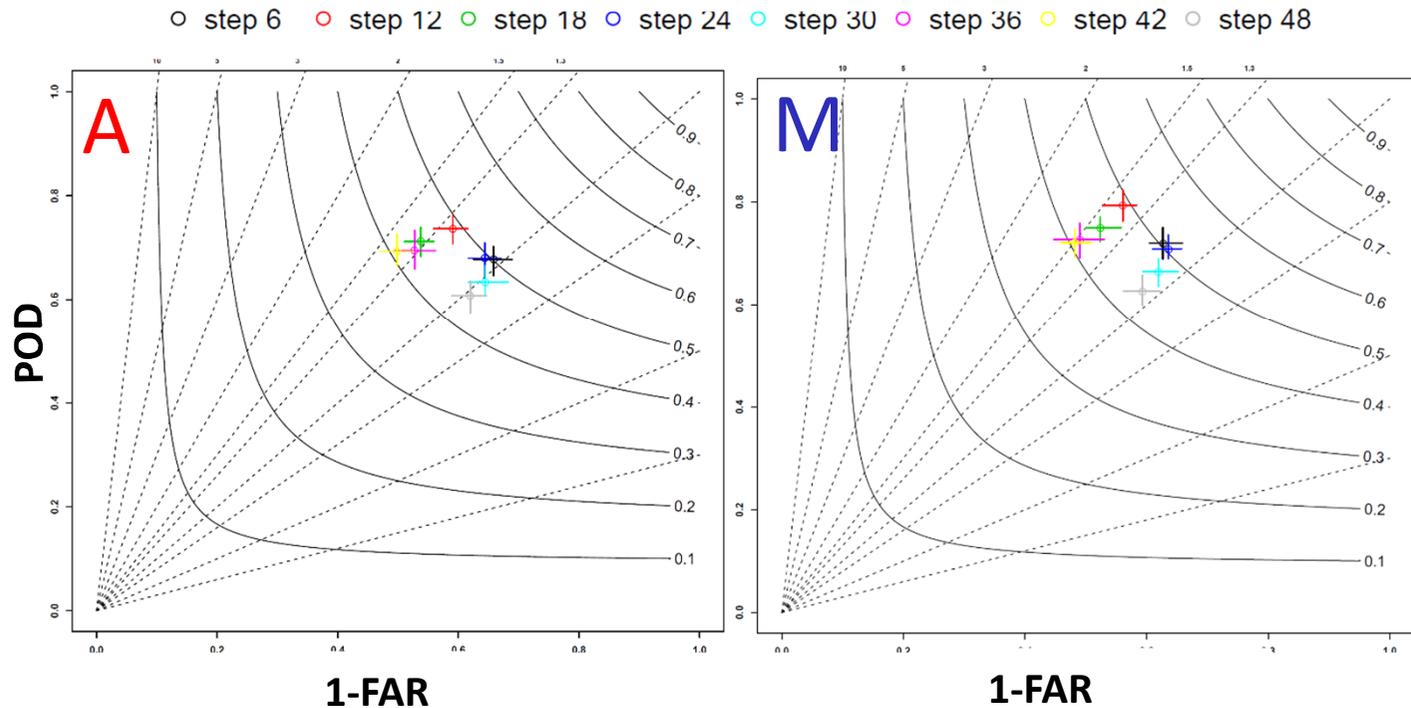
# Selection of an optimal advection scheme (MPDATA-A vs. MPDATA-M)



option **A** outperforms **M** for:  $T_{2m}$ , MSLP, 10-m wind speed,  
and for total cloud cover (the latter except only 36-48 h for RMSE)

# MPDATA-A vs. MPDATA-M

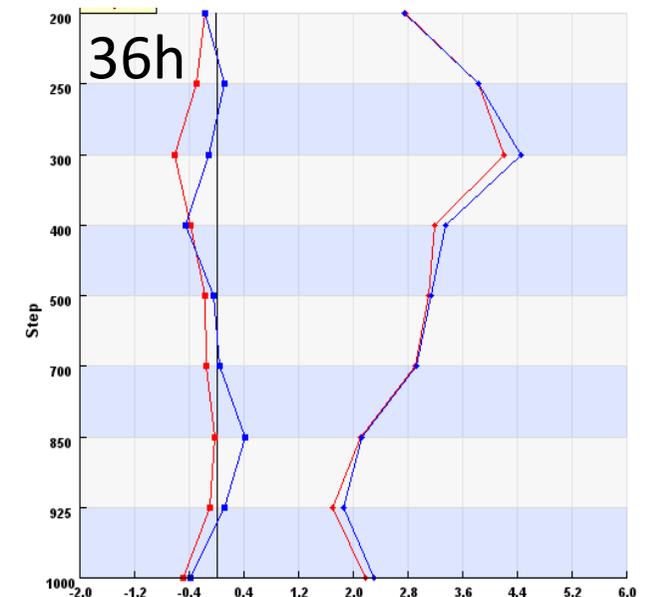
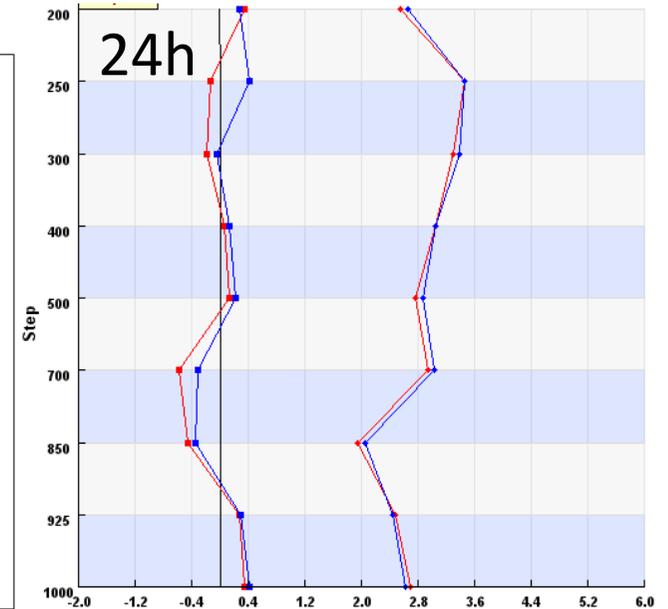
## Precipitation: 1mm and more



- For upper-air wind speed RMSE is usually lower for **A**
- **A** provides precipitation forecasts with slightly improved frequency bias
- Additionally: in the **A** simulations lower vertical velocities within convective updrafts are observed (not shown)

**COSMO-EULAG with the more diffusive scheme, MPDATA-A, provides forecasts having slightly better verification scores.**

## Upper-air wind speed



1. The more accurate MPDATA-M advection delivers worse scores than MPDATA-A. Possible reasons
  1. might be a hint for too less (horizontal) diffusion ?
  2. verification issue: better scores for more diffusive fields ?
2. Consistent, optimized and extensively tested COSMO-EULAG v5.5
3. The computational performance was slightly improved
4. COSMO-EULAG works semi-operationally in IMGW-PIB since winter 2019 with nudging and with competitive verification scores
5. Future work may involve comparison of COSMO-EULAG and ICON-LAM for high spatial resolutions (over Poland) and with more advanced verification

# A possible alternative dynamical core for ICON based on Discontinuous Galerkin Discretisation

Michael Baldauf (Deutscher Wetterdienst)

## MetStröm



## Discontinuous Galerkin (DG) methods in a nutshell

$$\frac{\partial q^{(k)}}{\partial t} + \nabla \cdot \mathbf{f}^{(k)}(q) = S^{(k)}(q), \quad k = 1, \dots, K$$

weak formulation  $\int_{\Omega_j} dx v(\mathbf{x}) \cdot \dots$

Finite-element ingredient

$$q^{(k)}(x, t) = \sum_{l=0}^p q_{j,l}^{(k)}(t) p_l(x - x_j)$$

e.g. Legendre-Polynomials

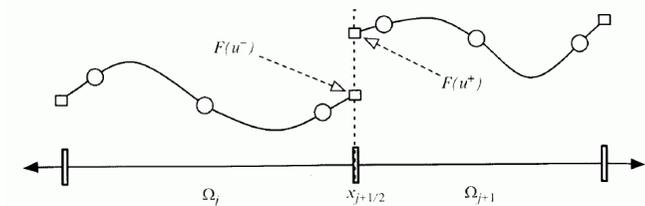
Finite-volume ingredient

$$\mathbf{f}(q) \rightarrow \mathbf{f}^{num}(q^+, q^-) = \frac{1}{2} (\mathbf{f}(q^+) + \mathbf{f}(q^-) - \alpha(q^+ - q^-))$$

Lax-Friedrichs flux

Gaussian quadrature for the integrals of the weak formulation

→ ODE-system for  $q^{(k)}_{jl}$



From *Nair et al. (2011)* in  
'Numerical techniques for global atm.  
models'

e.g.  
*Cockburn, Shu (1989) Math. Comput.*  
*Cockburn et al. (1989) JCP*  
*Hesthaven, Warburton (2008):*  
*Nodal DG Methods*

- **local conservation**
- any **order of convergence** possible
- flexible application on **unstructured grids** (also dynamic adaptation is possible, h-/p-adaptivity)
- very good **scalability**
- **explicit** schemes are easy to build and are quite well understood
- higher accuracy helps to **avoid several awkward approaches** of standard 2<sup>nd</sup> order schemes: staggered grids (on triangles/hexagons, vertically heavily stretched), numerical hydrostatic balancing, grid imprints by pentagon points or along cubed sphere lines, ...

- high computational costs due to
  - (apparently) **small Courant numbers**
  - higher number of DOFs
- **well-balancing** (hydrostatic, perhaps also geostrophic?) in Euler equations is an issue (can be solved!)
- basically ,only‘ an **A-grid-method**, however, the ,spurious pressure mode‘ is very selectively damped!

## Target system: ICON model

(Zängl et al. (2015) QJRMS)

- operational at DWD since Jan. 2015 (global (13km) and nest over Europe (6.5km))
- convection-permitting (2.2km): Q4/2020
- horiz.: icosahedral triangle C-grid, vertic.: Lorenz-grid
- non-hydrostatic, compressible
- mixed finite-volume / finite-difference (mass, tracer mass conservation)
- predictor-corrector time-integration → overall 2<sup>nd</sup> order discretization

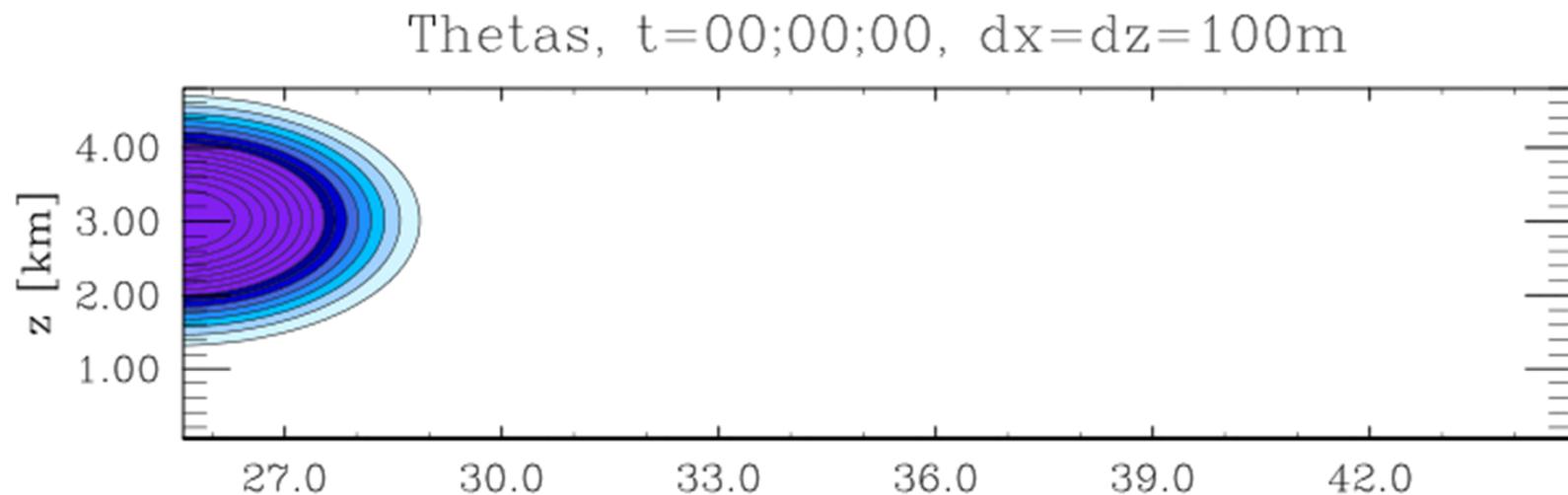


but currently far away from this, only a **toy model for 2D problems exists** with:

- explicit time integration DG-RK (with Runge-Kutta schemes) or horizontally explicit-vertically implicit (DG-HEVI) (with IMEX-Runge-Kutta)
- ‚local DG‘ (LDG) option for PDEs with higher spatial derivatives
- use of a triangle grid (also on the sphere) is optional

## Test case: falling cold bubble

Testsetup by *Straka et al (1993)*



Test properties:

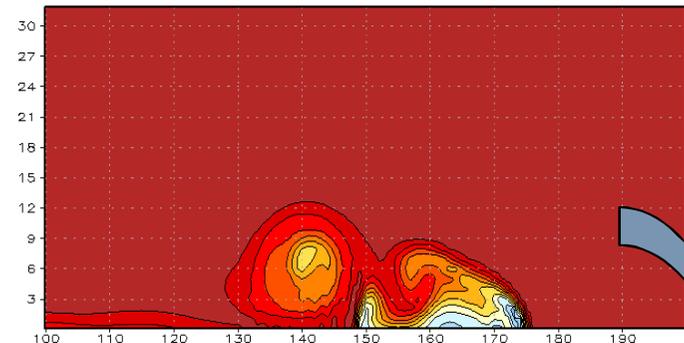
- test of dry Euler equations (without Coriolis force)
- unstationary
- strongly nonlinear
- comparison with reference solution from paper

## DG explicit

2<sup>nd</sup> order

dx=dz=200m

Theta p=1 t=900.0



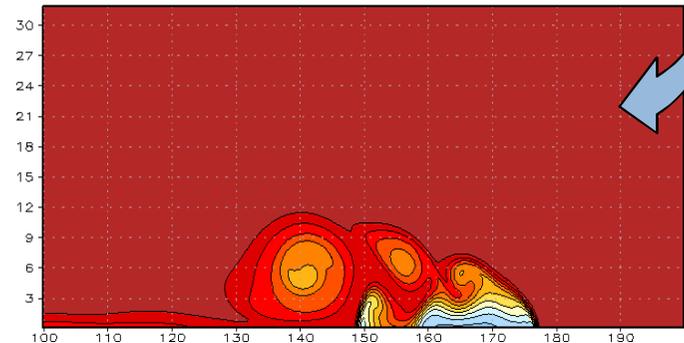
Theta: Mean: 299.593 Min: 290.014 Max: 300.361  
GrADS: COLA/IGES 2019-03-13-10:48

**Faktor 4.3**  
in comput. time

3<sup>rd</sup> order

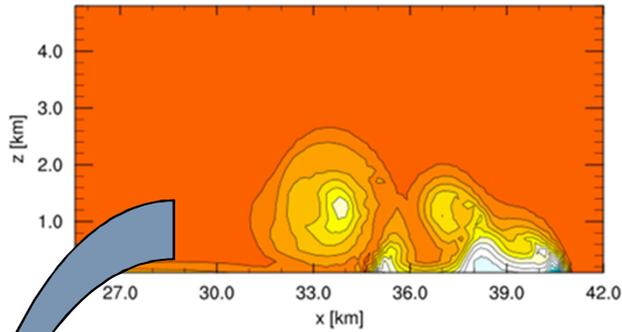
dx=dz=200m

Theta p=2 t=900.0

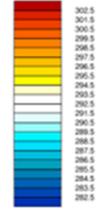


Theta: Mean: 299.595 Min: 290.303 Max: 300.295  
GrADS: COLA/IGES 2019-03-13-10:42

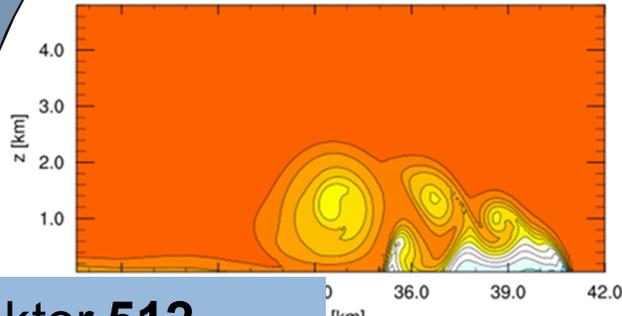
t=00:15:00, dx=dz=200m,



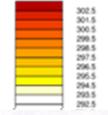
Theta (in K)



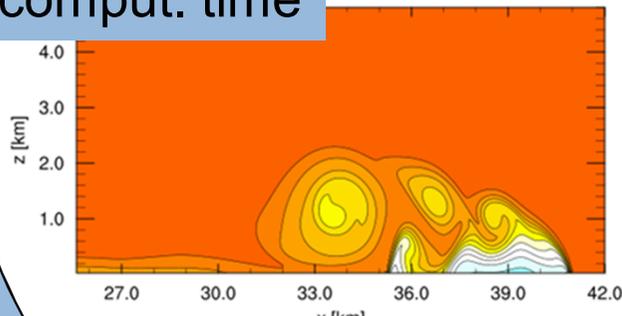
t=00:15:00, dx=dz=100m,



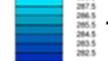
Theta (in K)



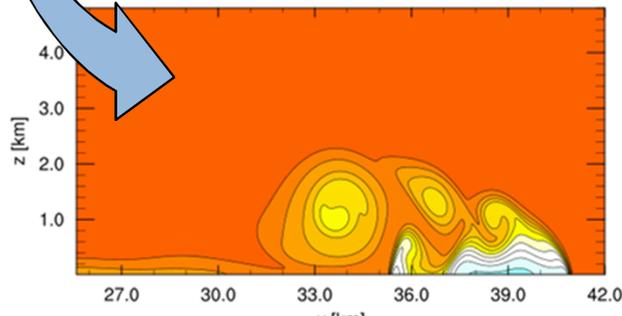
t=00:15:00, dx=dz=50m,



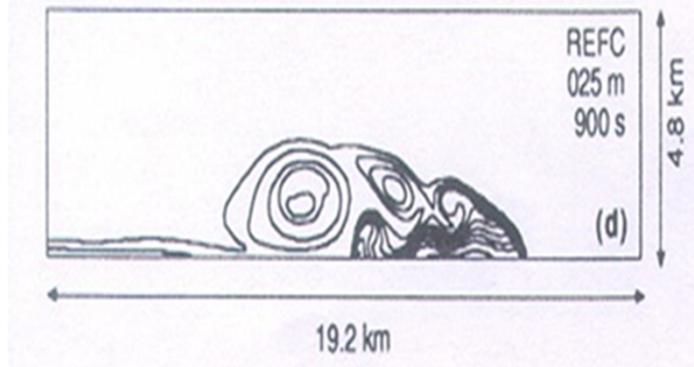
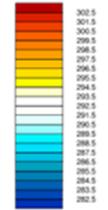
Theta (in K)



t=00:15:00, dx=dz=25m,



Theta (in K)



Reference solution  
from Straka et al. (1993)

**Faktor 512**  
in comput. time

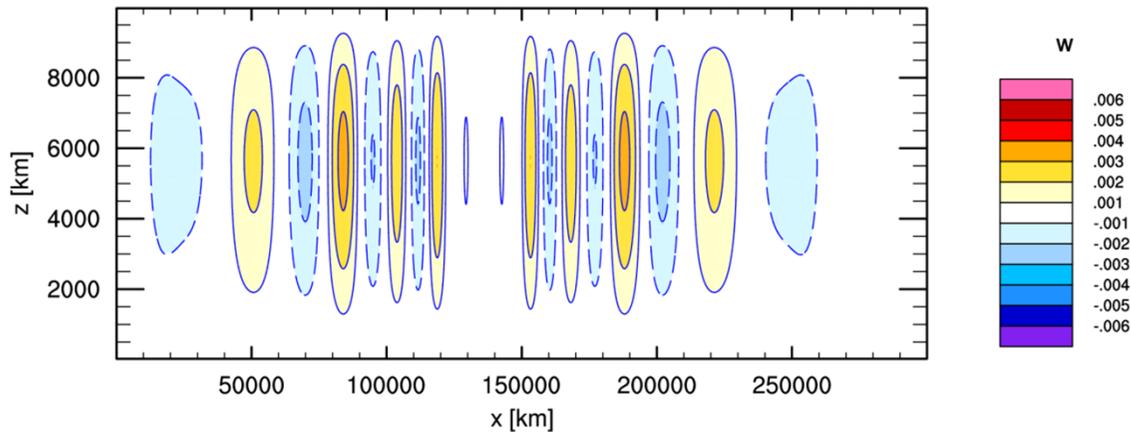
# Linear gravity/sound wave expansion in a channel



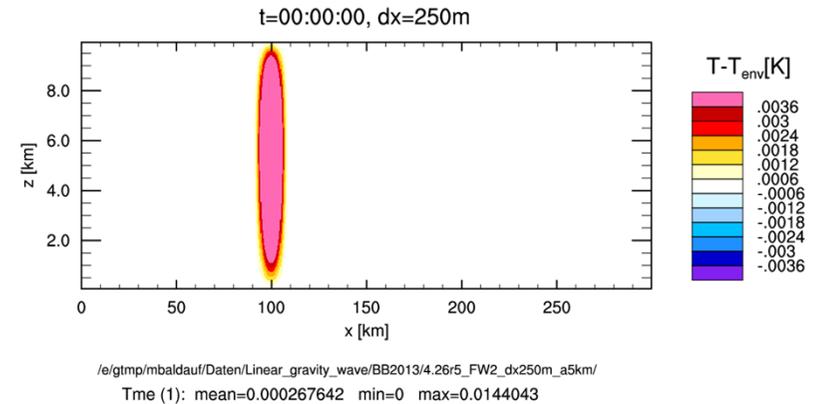
setup similar to *Skamarock, Klemp (1994) MWR*

$$\Delta x = 500\text{m}, \Delta z = 250\text{m}$$

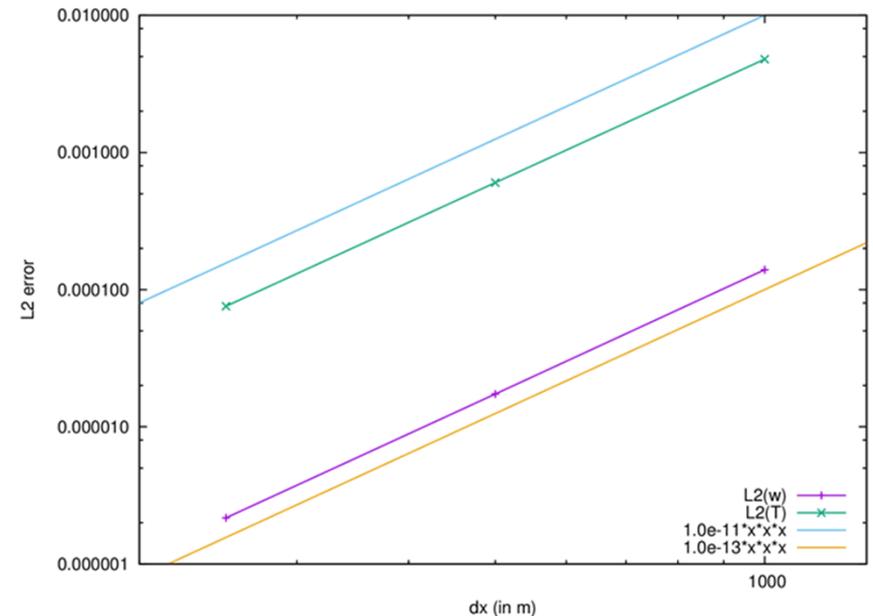
w, t=00:30:00



colors : simulation with p=2/RK3-SSP (i.e. 3<sup>rd</sup> order explicit DG)  
blue lines: analytic solution for compressible, non-hydrostatic Euler eqns. (*Baldauf, Brdar (2013) QJRMS*)



Exact 3<sup>rd</sup> order convergence for w and T':





## Horizontally explicit - vertically implicit (HEVI)-scheme with DG

*Motivation:* get rid of the strong time step restriction by vertical sound wave expansion in flat grid cells (in particular near the ground)

$$\frac{\partial q^{(s)}}{\partial t} + \underbrace{\nabla \cdot \mathbf{f}_{slow}^{(s)}}_{\text{explicit}} + \underbrace{\nabla \cdot \mathbf{f}_{fast}^{(s)}}_{\text{implicit}} = \underbrace{S_{slow}^{(s)}}_{\text{explicit}} + \underbrace{S_{fast}^{(s)}}_{\text{implicit}}$$

$$\mathbf{f}_{fast}^{(s)} = f_{z,fast}^{(s)} \mathbf{e}_z$$

$$f_{z,fast}^{(s)} = \sum_{s'} H^{ss'} q^{(s')}$$

- Use of IMEX-RK (SDIRK) schemes: SSP3(3,3,2), SSP3(4,3,3) (*Pareschi, Russo (2005) JSC*)
- The implicit part leads to several band diagonal matrices  
→ here a direct solver is used (expensive!)

References:

*Giraldo et al. (2010) SIAM JSC:* propose a HEVI semi-implicit scheme

*Bao, Klöfkorn, Nair (2015) MWR:* use of an iterative solver for HEVI-DG

*Blaise et al. (2016) IJNMF:* use of IMEX-RK schemes in HEVI-DG

*Abdi et al. (2017) arXiv:* use of multi-step or multi-stage IMEX for HEVI-DG

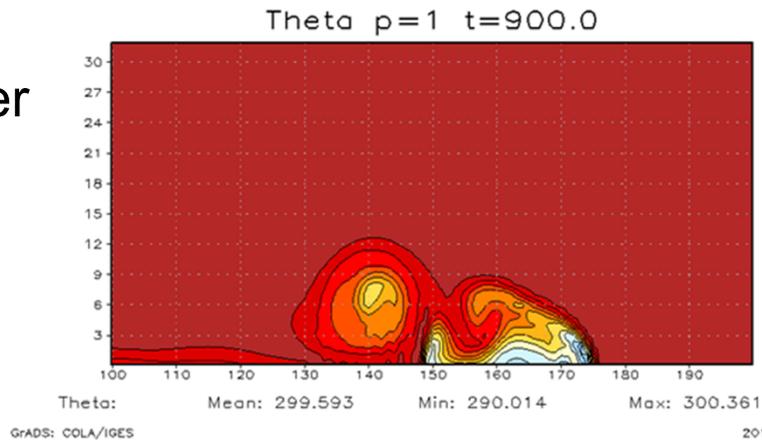
# Test case: falling cold bubble (Straka et al. (1993))

## Comparison explicit vs. HEVI scheme

2<sup>nd</sup> order

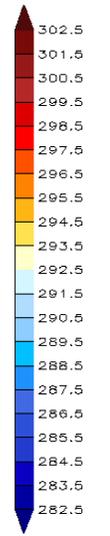
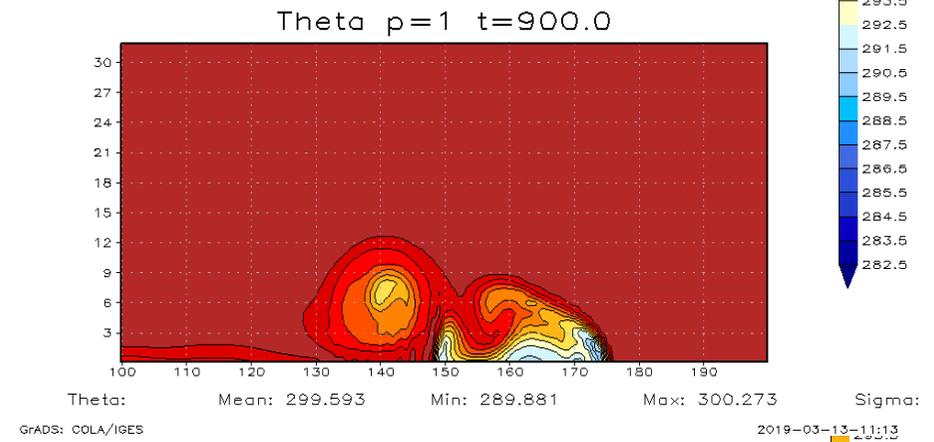
RK2-SLC  
dt=0.0826446280992  
dx=200.0  
dy=200.0

DG explicit

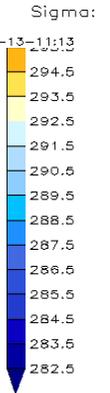
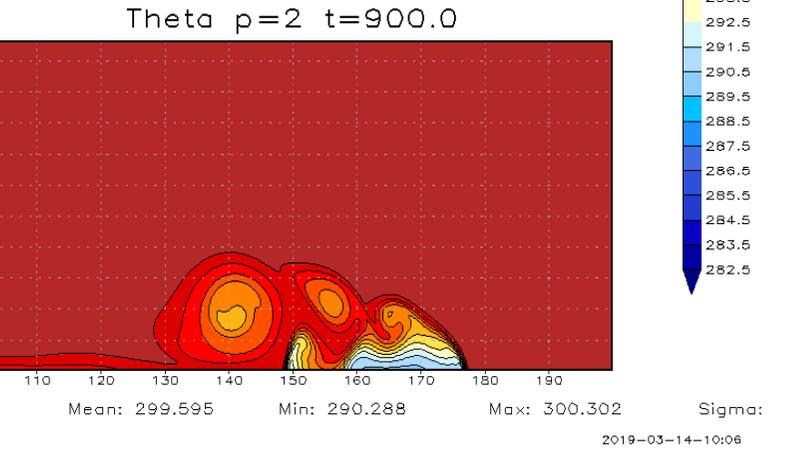
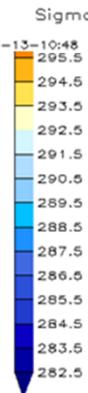
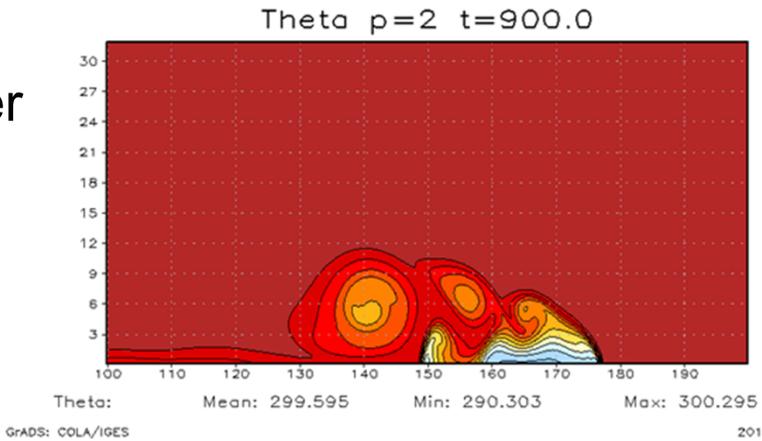


SSP3-3-3-2  
dt=0.16393442623  
dx=200.0  
dy=200.0

DG HEVI



3<sup>rd</sup> order



## How to bring DG on the sphere ...

Idea to avoid pole problem and to keep high order discretization:  
use **local (rotated) coordinates** for every (triangle) grid cell,  
i.e. rotate every grid cell towards  $\lambda \approx 0, \varphi \approx 0$ .

→ geometry is treated exactly

→ transform fluxes between neighbouring cells

shallow water equations

covariant formulation (here: without bathymetry)

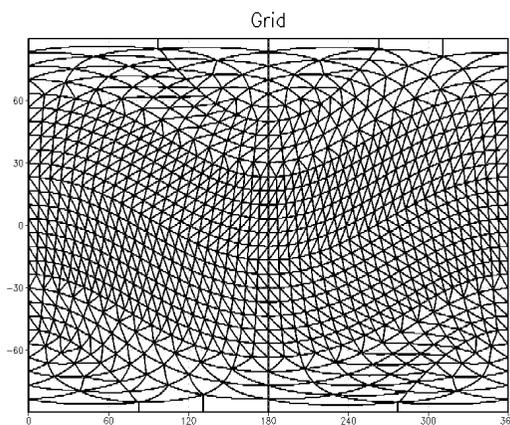
$$\begin{aligned} \frac{\partial \sqrt{GH}}{\partial t} + \frac{\partial}{\partial x^i} \sqrt{G} m^i &= 0 \\ \frac{\partial \sqrt{G} m^i}{\partial t} + \frac{\partial}{\partial x^j} \sqrt{G} T^{ij} &= \sqrt{G} (F_{Cor}^i - \Gamma_{jk}^i T^{jk}) \\ T^{ij} &= \frac{m^i m^j}{H} + \frac{1}{2} g^{ij} g_{grav} H^2 \end{aligned}$$

# Barotropic instability test

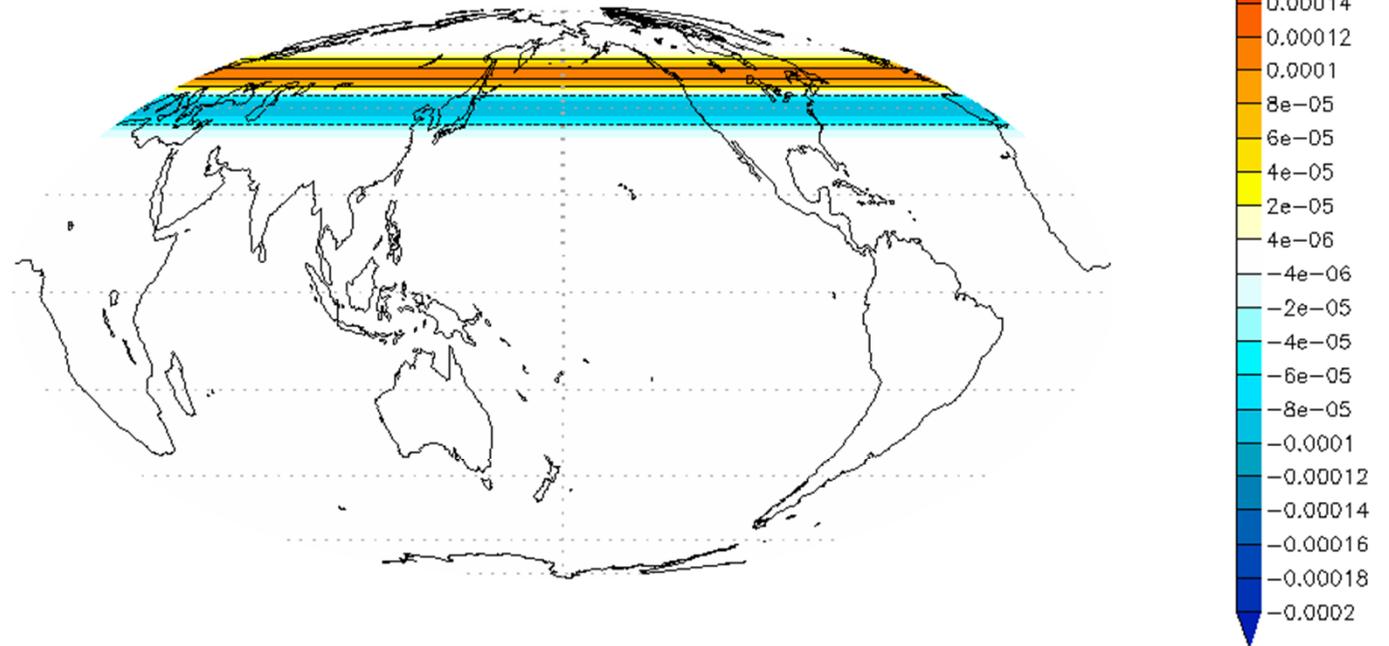
*Galewsky et al. (2004)*

**4th order DG scheme**  
without additional diffusion  
dx~67 km, dt=15 sec.

simple triangle grid  
on the sphere  
dx ~ 500km:



rel. Vortic., ord=4 t=0d00h00m0.0s



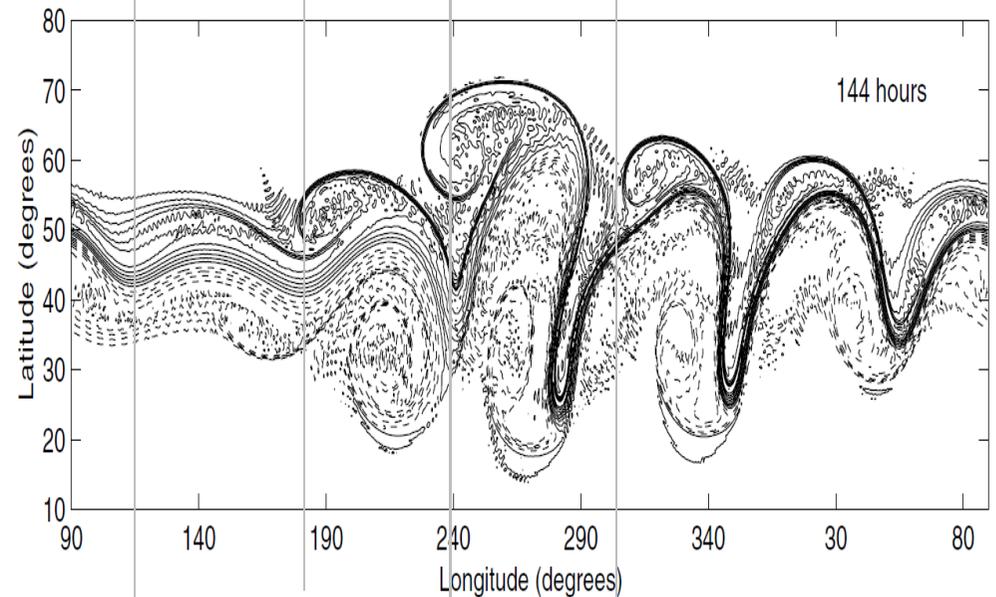
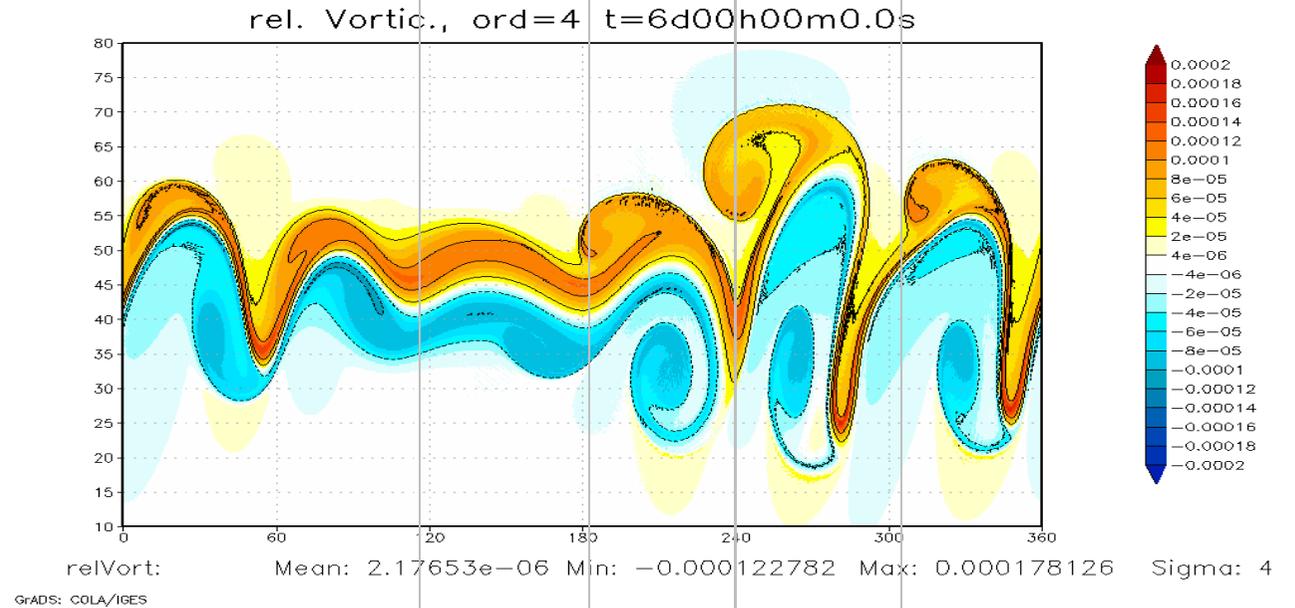
# Barotropic instability test *Galewsky et al. (2004)*

**4th order DG scheme**  
without additional diffusion  
dx~67 km, dt=15 sec.

**FMS-SWM** (Geophys. Fl. Dyn. Lab.)  
without additional diffusion  
dx~60 km (T341), dt=30 sec.

Fig. 4 from *Galewsky et al. (2004)*

## relative vorticity

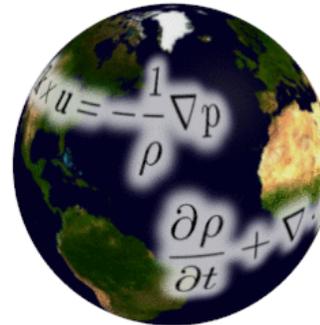


## Summary

- **2D toy model** for
  - explicit DG-RK (on arbitrary unstructured grids with triangle or quadrilateral grid cells) and
  - **HEVI DG-IMEX-RK**works for several idealized tests (also Euler equations with terrain-following coordinates), correct convergence behaviour, ...
- **DG on the sphere** by use of local (rotated gnomonial) coordinates

## Outlook

- **further design decisions:** nodal vs. modal, local DG vs. interior penalty vs. .... ..
- coupling of **tracer advection** (mass-consistency)?
- improve **efficiency** in the HEVI direct solver
- further **milestones** (for the next years!)
  - development of a 3D prototype DG-HEVI solver
  - choose optimal convergence order  $p$  and grid spacing  
estimated:  $p_{\text{horiz}} \sim 3 \dots 6$ ,  $p_{\text{vert}} \sim 3 \dots 4$  ( $p_{\text{time}} \sim 3 \dots 4$ )



PDEs

THE WORKSHOP ON PARTIAL  
DIFFERENTIAL EQUATIONS ON THE SPHERE

Announcement:

The next

**„Partial differential equations on the sphere“ – workshop**

will take place at  
DWD, Offenbach, Germany  
5-9 October 2020