

High order conserving schemes on arbitrary polygonal grids

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Bad Orb 2009

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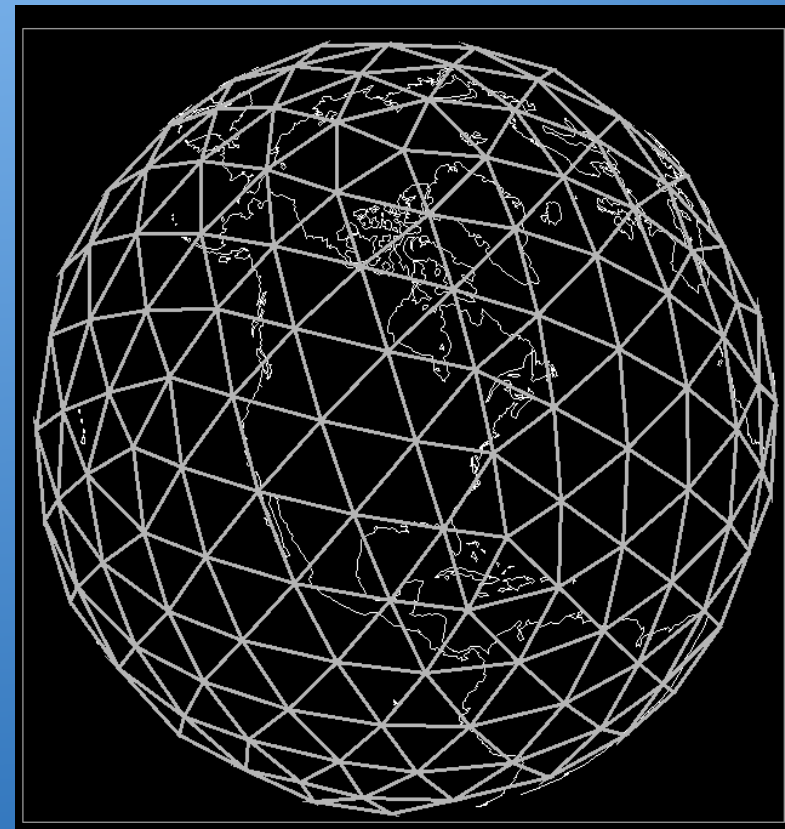
Conservation and high order

- \top The Importance of Conservation
- The Importance of Maintaining Second or Third Order Accuracy:
 1. Bulls eyes at special points of the icosahedral grid
 2. Big problems with sudden grid changes in 1-d
 - Many models Use Grid Regularisation on the Icosahedron and a Smooth Vertical Grid Structure
 - Problems with Irregular Resolution can be Traced Back to First Order approximations at Special Points
 - Schemes with Uniformly Second or Third Order Approximation Have no Bulls Eyes and are Robust

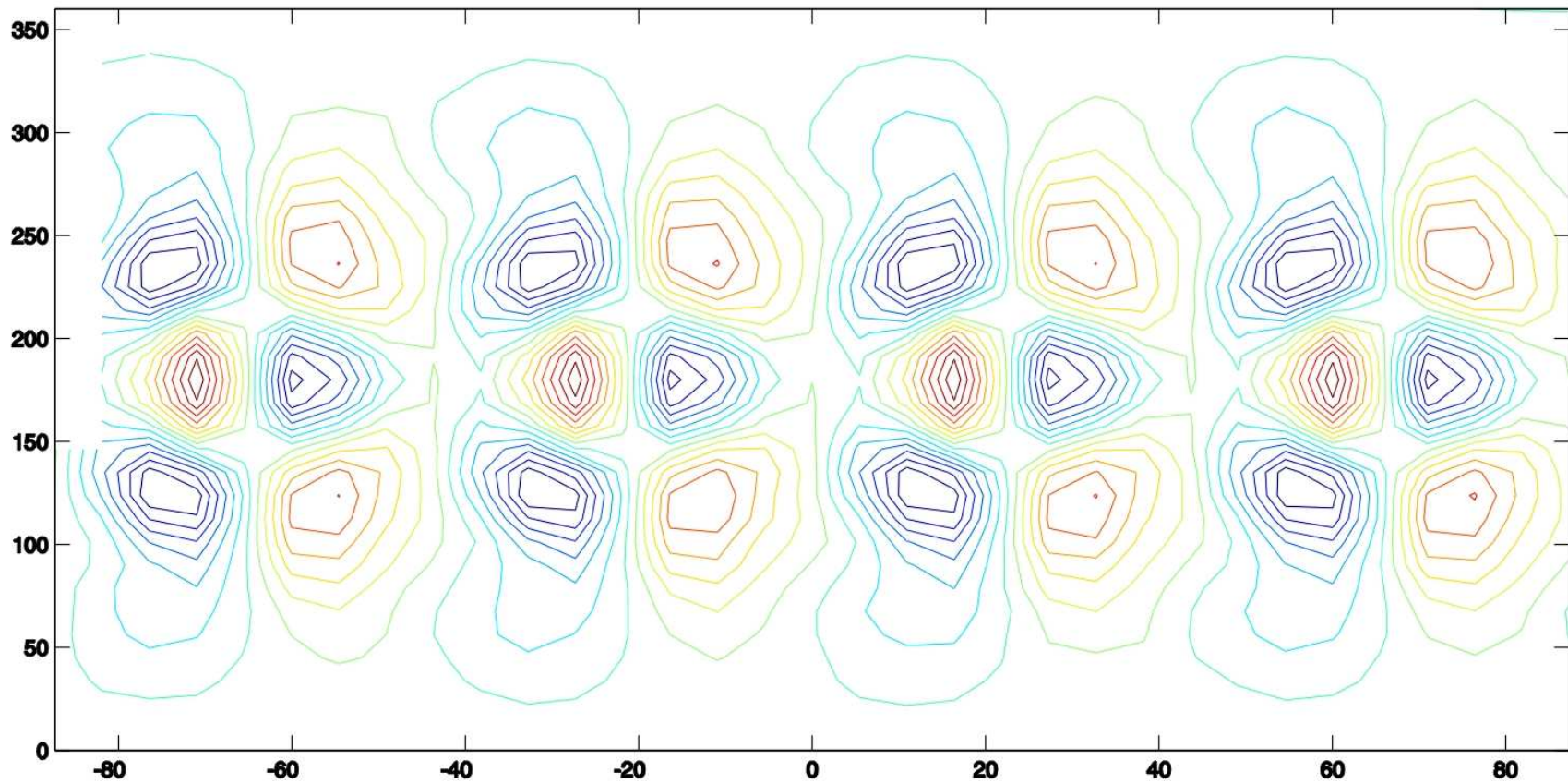
Icosahedral Grids

Rhomboidal Grid: hexagonal and pentagonal stencils and dual cells

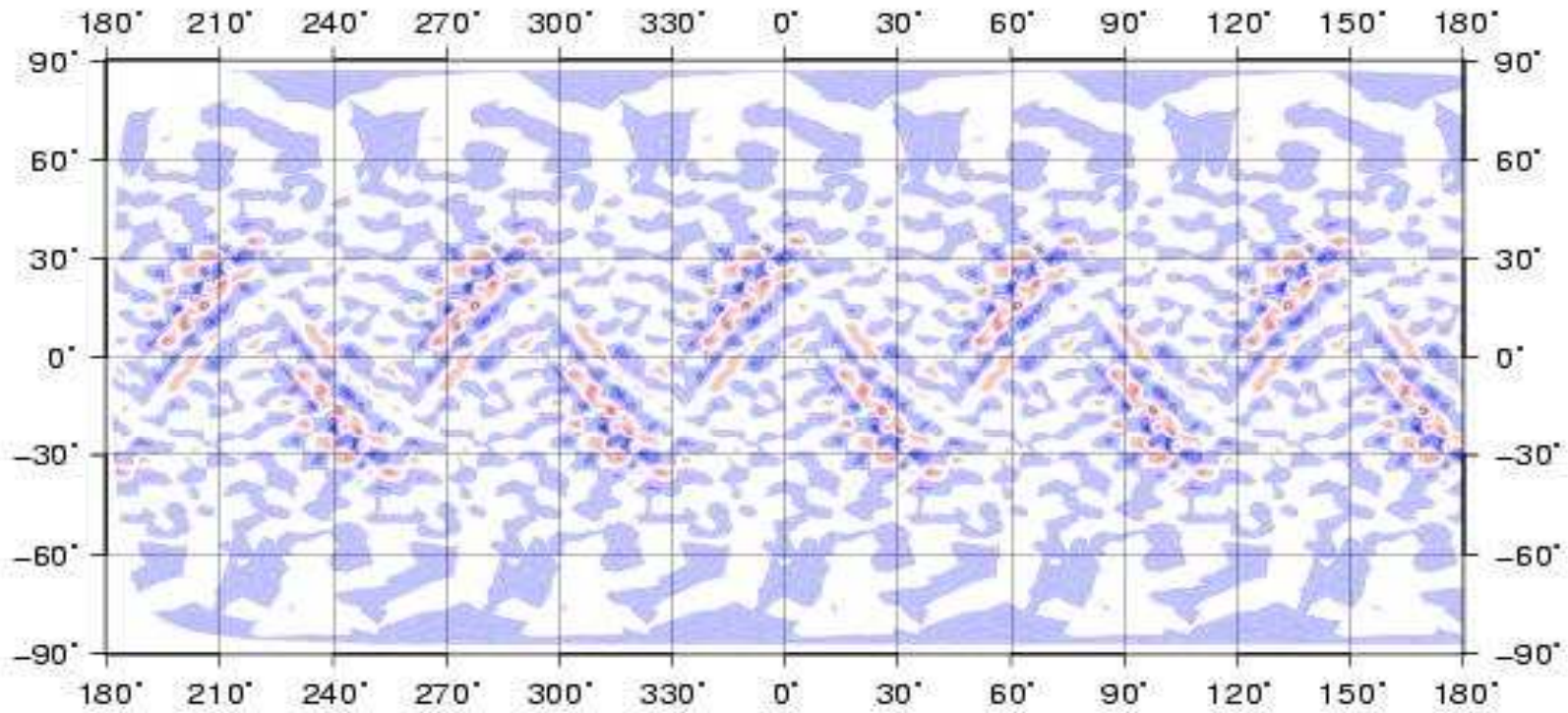
- An Uniformly Third Order Scheme was implemented and produces no Bulls Eyes (MWR 136 (2008) pp 2483, Steppeler and Ripodas)
- Are ***flux based conserving second or third order schemes*** possible ?
(YES)



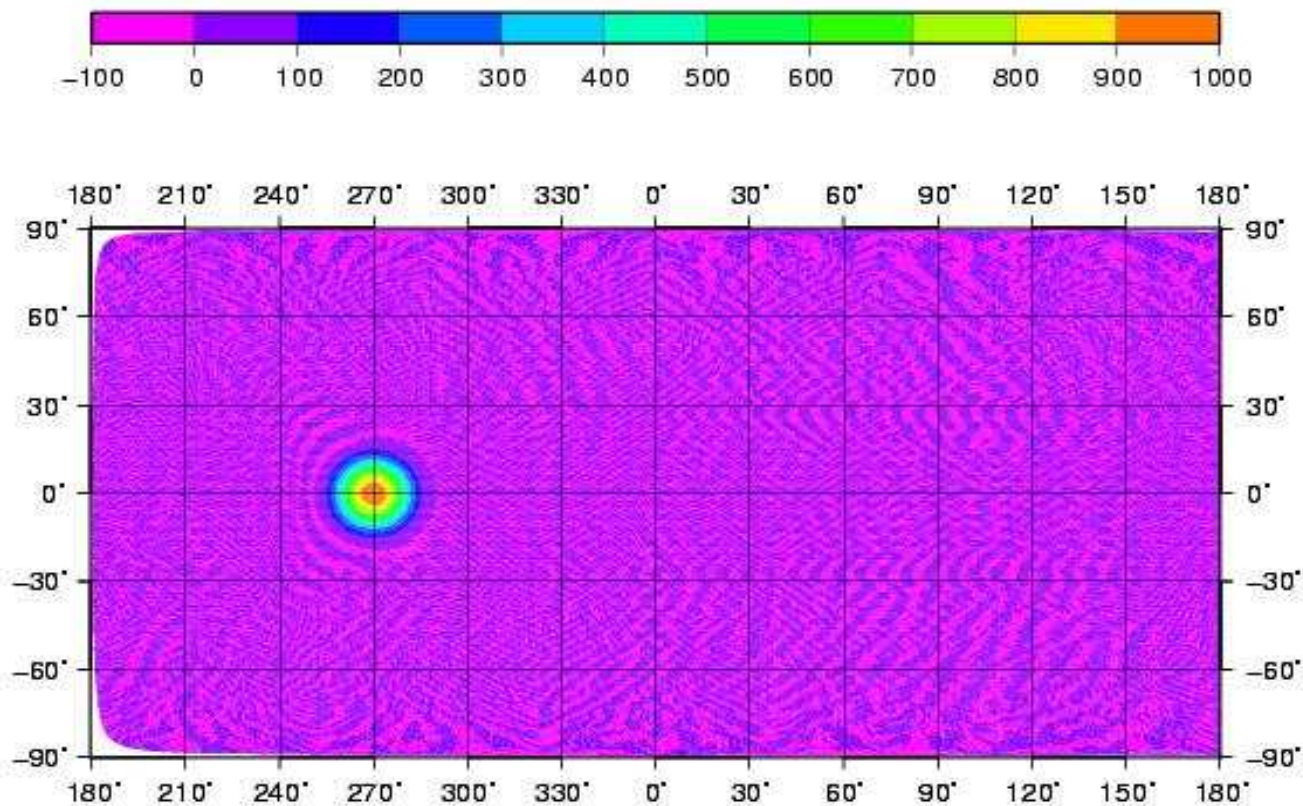
No Diffusion, Forecasted Divergence After 2,5h



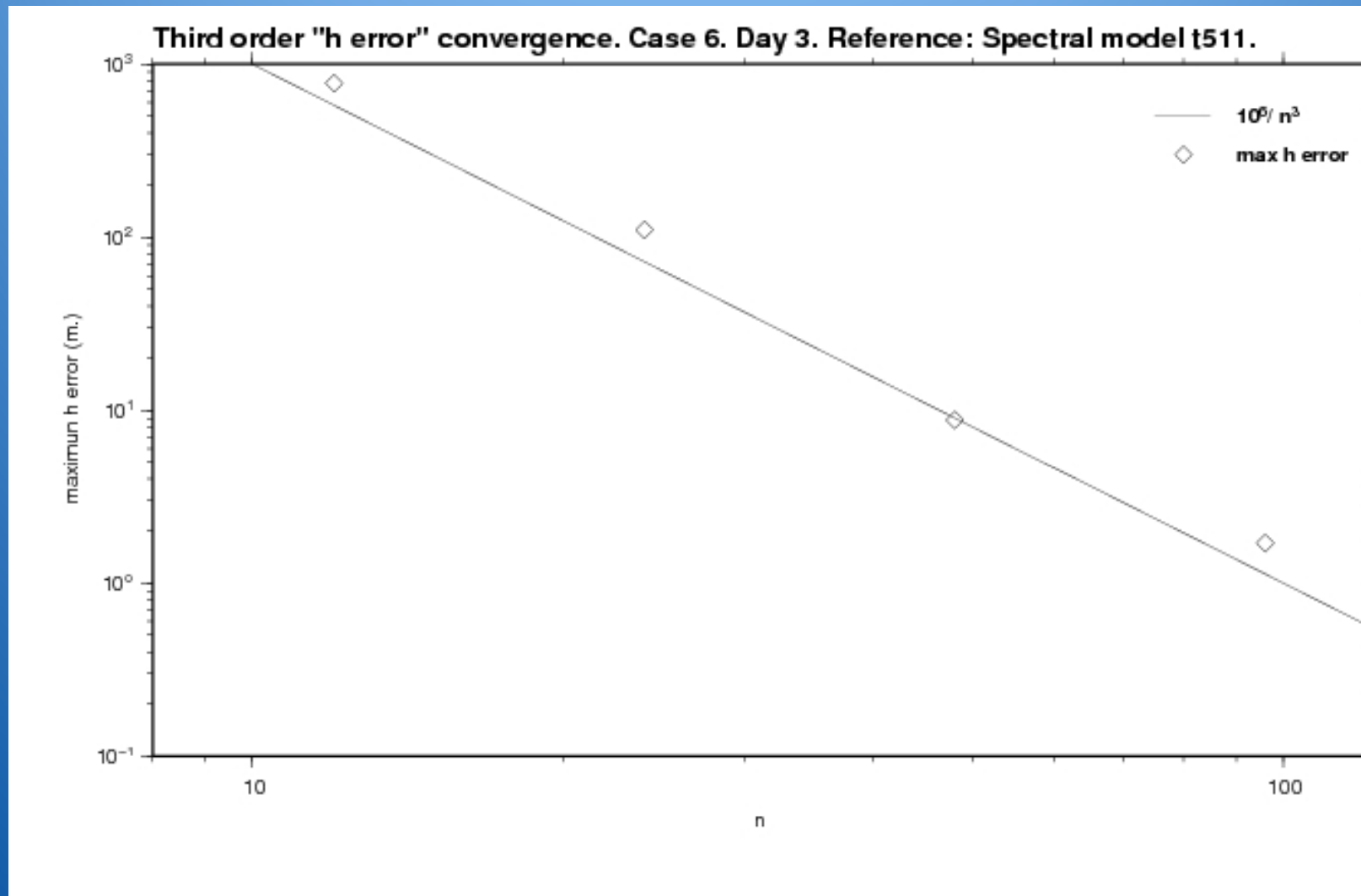
Solid Body Rotation



Homogeneous Advection



Third Order Convergence of Shallow Water Model at Day 3



Grid Stencils Baumgardner

- Edges grid

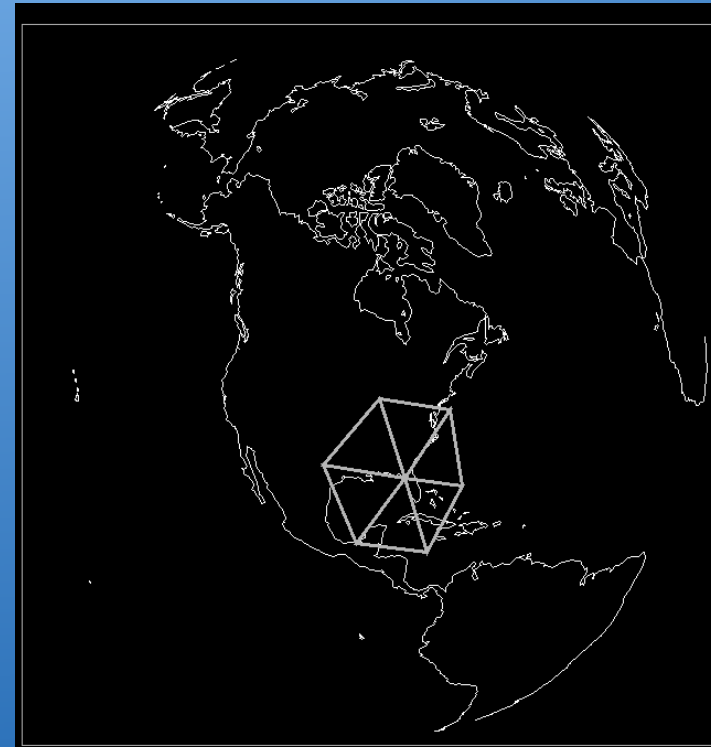
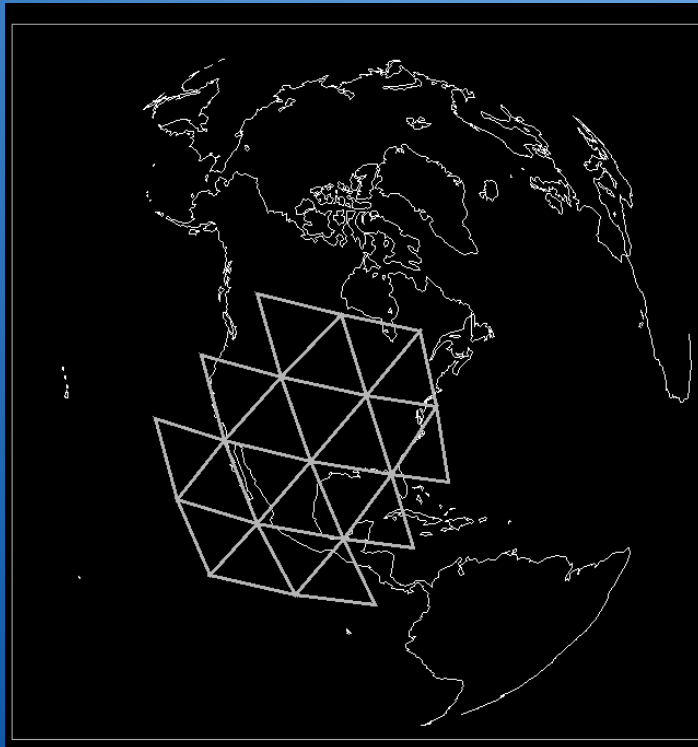
- Order 3

Redundancy 19:10

- Edges grid

- Order 2

Redundancy 7:6 or 6:6



High and low order mass formula

Spectral, spectral elements, finite elements:

$$mass_element = \int_{dF} spectral_rep_of_density \circ dx$$

Finite Difference:

$$mass_element = \int_{dF} \rho dx \approx \rho_i dx$$

Question: is FD conservation possible with a high order approximation of the density equation?

Definition of flux based conservation schemes

$$\frac{\partial h}{\partial t} = \frac{\partial uh}{\partial x}; f = uh; u = 1$$

$$\frac{\partial h}{\partial t} = \frac{\partial h}{\partial x}; f = h$$



Finite difference flux based conservation:

$$dx \frac{\partial h_0}{\partial t} = h_{1/2} - h_{-1/2}; \Rightarrow \frac{\partial H}{\partial t} = \sum_i dx \frac{\partial h_i}{\partial t} = 0$$

How to define point fluxes $f^*_{-1/2}, f^*_{1/2}$

Stencil: $i \in \{-2, -1, 0, 1, 2\}$

Point flux definition: $f^*_{-1/2} = \sum_i a_i f_i; f^*_{1/2} = \sum_i b_i f_i$

Equation for the a_i, b_i : $(f^*_{1/2} - f^*_{-1/2}) / dx - \frac{\partial f}{\partial x} = o(dx^5)$

for all polynomial functions $f(x) = x^n, n \in \{0, 1, 2, 3, \dots\}$

Simple example in 1-d

(Only for **regular** grids or the high resolution limit)

$$fp_{i+1} = f_{i+1} \frac{4}{3} - (f_{i+2} + f_i) \frac{1}{6};$$
$$fp_{i-1} = f_{i-1} \frac{4}{3} - (f_{i-2} + f_i) \frac{1}{6}$$

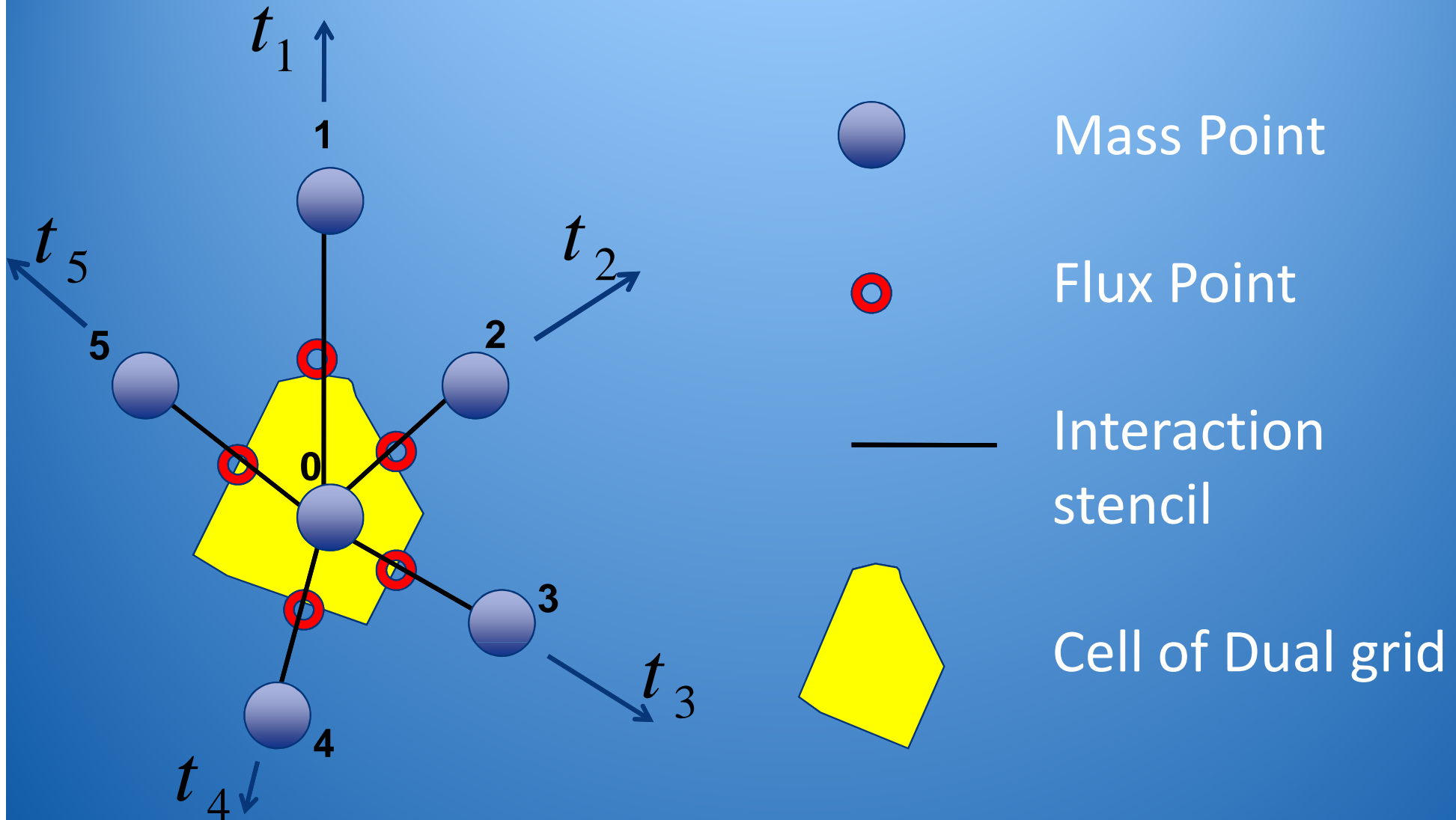
$$\frac{\partial f}{\partial x} 2 dx = fp_{i+1} - fp_{i-1} = \frac{4}{3} (f_{i+1} - f_{i-1})$$
$$- \frac{1}{3} (f_{i+2} - f_{i-2}) / 2 + o(x^5)$$

Properties of flux based conservation schemes

- If the $h_{1/2}$ are exactly computed by spectral method, the flux based scheme is of order 2 at most
- If the $h_{1/2}$ are computed by fourth order interpolation from the h_i , the flux based scheme is of order 2 at most
- **If $h_{1/2}$ deviate from the exact value by more than order 2, it is possible to define them in such a way to result into a 3rd or 4th order approximation of the density equation.** Such coarsely approximated fluxes are called **point fluxes.**

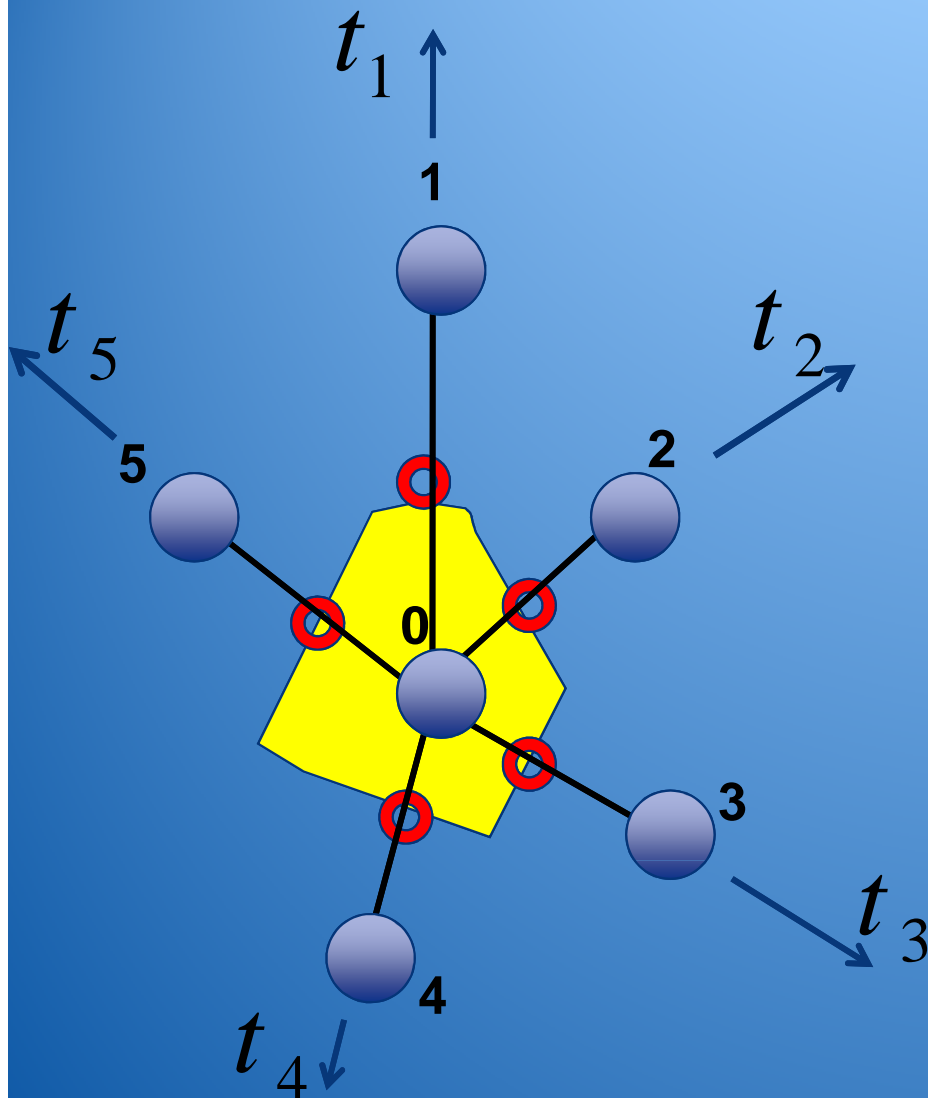
- The example generalises to the regular square and hexagon
- To derive similar point fluxes for irregular grids is not trivial
- The formula for the irregular grid to be given are **not** necessarily identical to that above when specialised to the regular case
- **The grid regularisation procedures (spring dynamic) are not necessary when using the Point fluxes for third or second order for the irregular grid**

Irregular Pentagonal Cell



The grid is called quasi regular, if stencil lines cut the sides of the dual cell in half

Irregular Pentagonal Cell



t_1 : directional parameter
for line 0,1 etc

f_1 : point flux for point
on line 0,1 etc

ρ_1 : density derivative for
point 0 etc

w_1 : weight to be defined
for the point flux

$\frac{\delta f_i}{\delta t_i}$: directional derivative

The definition of weights and point fluxes will be obtained by the geometric and collocation formula . This is non trivial even in 1-d for irregular resolution.

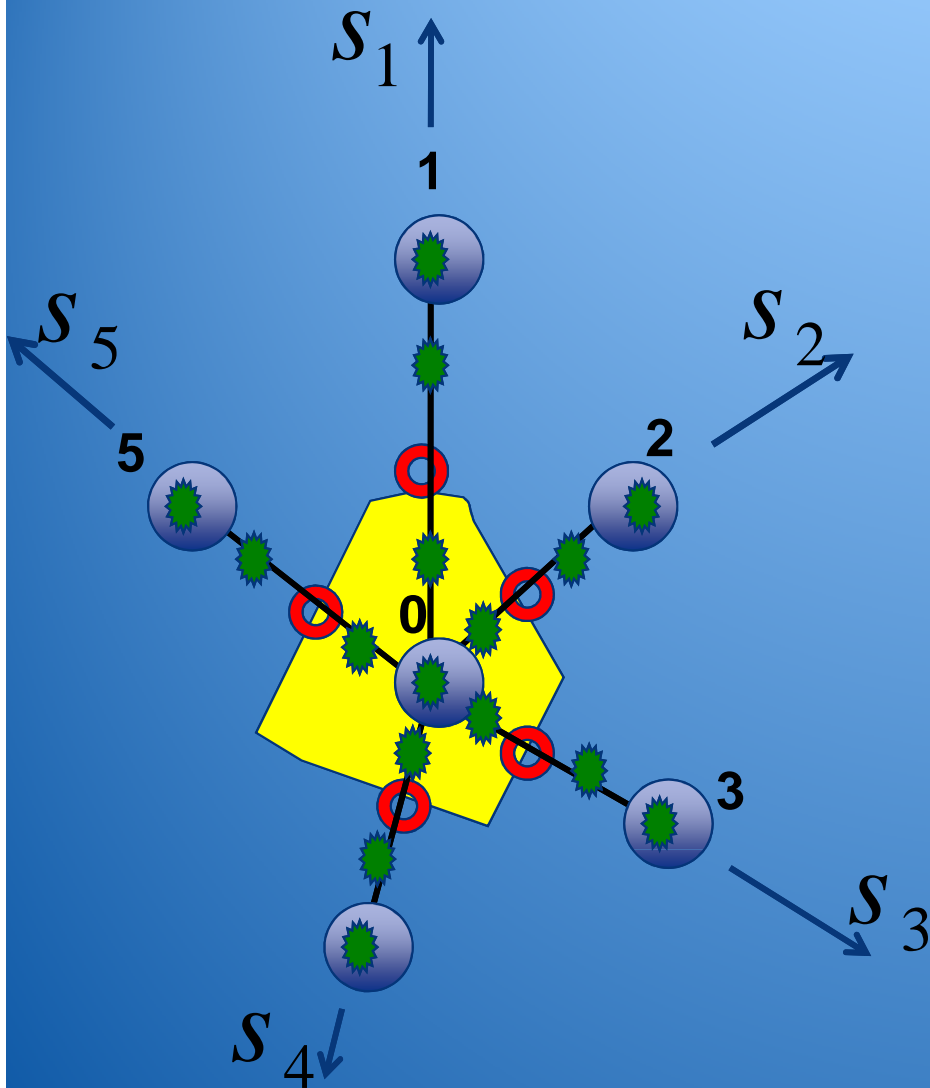
Irregular Pentagonal Cell

Geometric Formula,
derived from Stokes Eq:

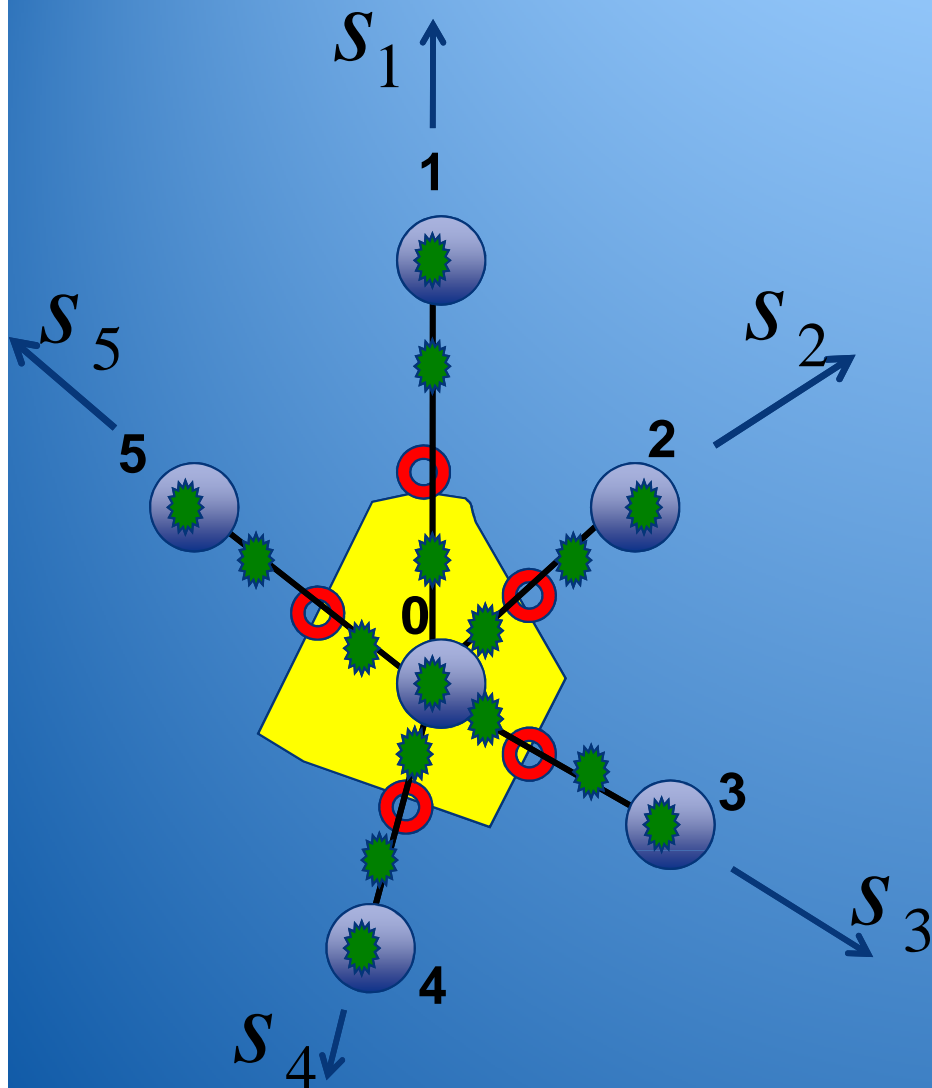
$$\frac{\partial \rho}{\partial t} = \sum_i w_i = \sum_i w_i \vec{f}_i \cdot \frac{\delta f}{\delta s_i} (P_0) + w_i iht_i$$

$$iht = \frac{\delta f}{\delta s_i^\perp} (P_0)$$

The inhomogeneous term iht is different from 0, when the grid is not quasi regular



Irregular Pentagonal Cell

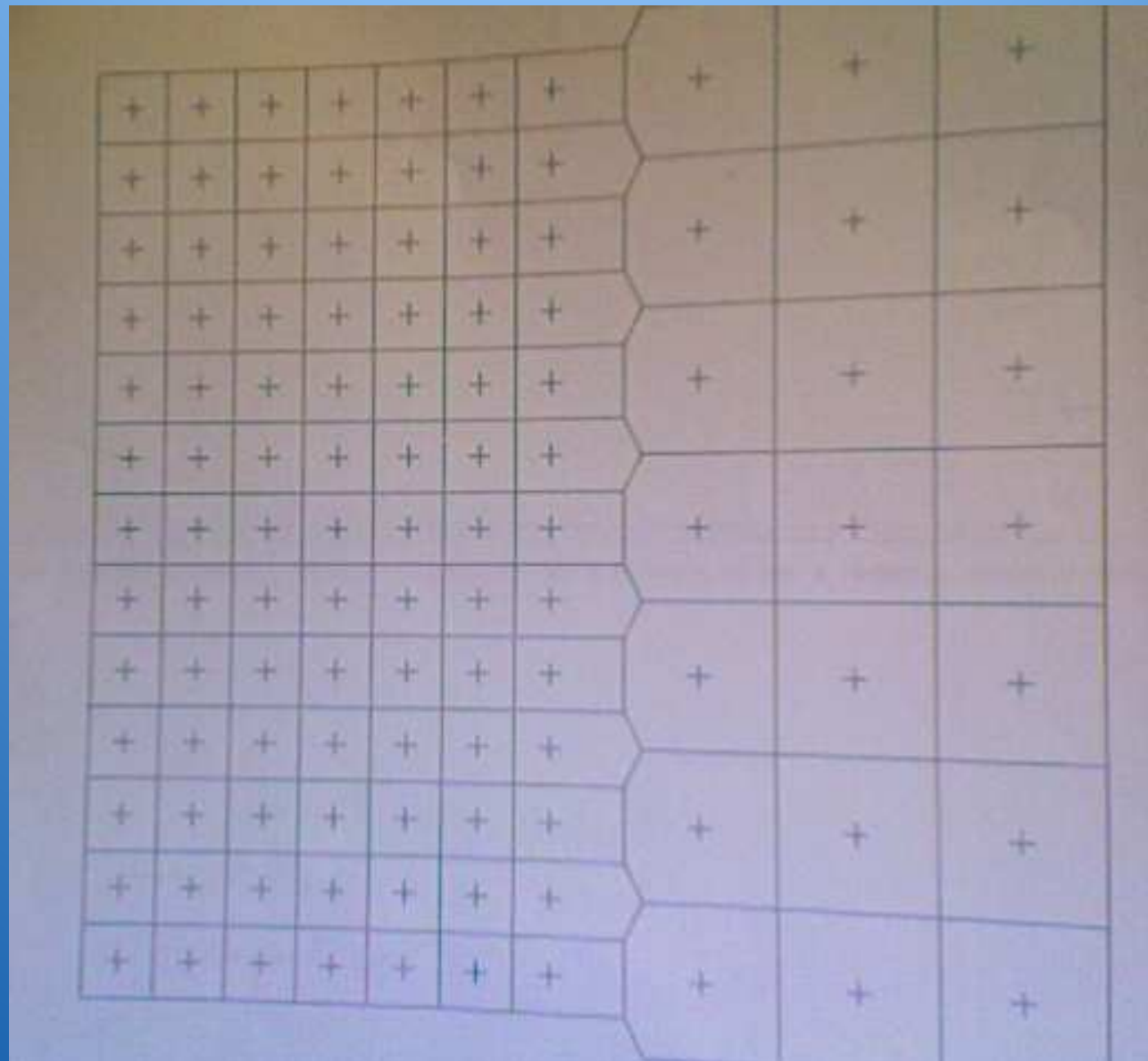


The scheme defined above is conserving, if the point fluxes defined above for points 0 and 1 etc compensate. This is achieved by defining the derivatives by collocation points obtained by interpolation or Galerkin methods.

1. In an irregular grid even order 2 in 1-d is non trivial
2. It is a generalisation of Arakawa method, but all interpolations and differentiations must be done in third (second) Order.

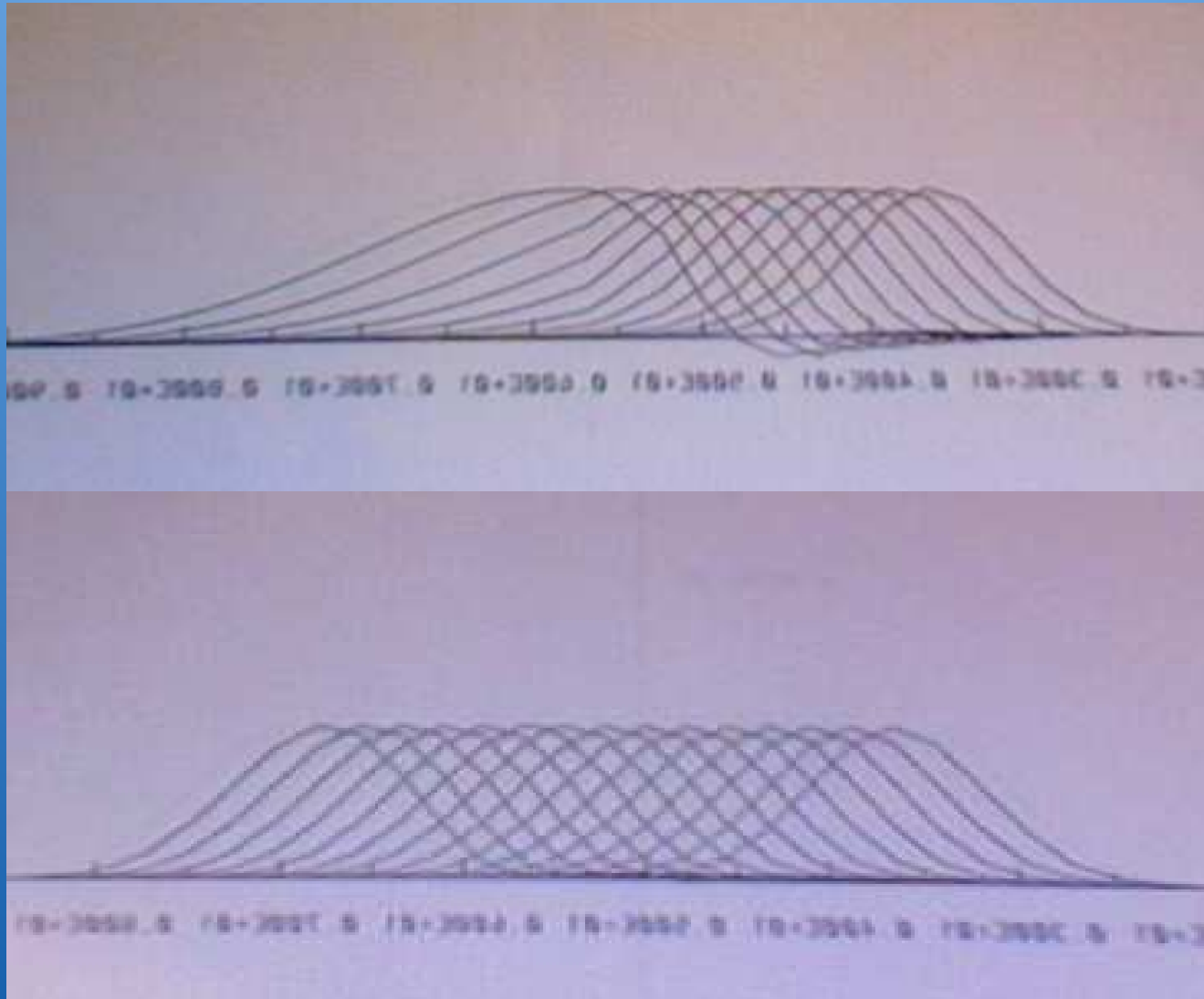
Dual Cells with Pentagons

2-d Example: RESOLUTION CHANGE USING



1-d Example

3-rd order scheme with conservation



Conclusions

- In an irregular cell structure point fluxes can be defined as linear functions of surrounding mass or flux points, leading to conserving second or third order schemes.
- Each point flux uses only a small part of the (large) Baumgardner stencil.
- Point fluxes cannot be accurate approximations of fluxes, such as high order interpolations of grid values of fluxes.
- The schemes are free of grid generated noise (Bulls eyes on irregular grids).

Conclusions (Cnt)

- The definition of point fluxes is not unique, therefore other restraints, such as desired behaviour of $\text{div } 0$ waves, can be investigated
- The schemes can be implemented easily, when coordinate of points and the neighbourhood table for mass and flux points as well as the neighbourhood information (Stencil) is available.
- Cut cells can be brought from first to higher order (for the cells near boundaries).
- Problems encountered with irregular and adaptive resolution can often be traced back to discretisation crimes, even hanging nodes.